

# First principle approach to correlation functions of spin-1/2 Heisenberg chain : fourth-neighbor correlators

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## Abstract

We show how correlation functions of the spin-1/2 Heisenberg chain without magnetic field in the anti-ferromagnetic ground state can be explicitly calculated using information contained in the quantum Knizhnik-Zamolodchikov equation [qKZ]. We find several fundamental relations which the inhomogeneous correlations should fulfill. On the other hand, it turns out that these relations can fix the form of the correlations uniquely. Actually, applying this idea, we have obtained all the correlation functions on five sites. Particularly by taking the homogeneous limit, we have got the analytic form of the fourth-neighbor pair correlator  $\langle S_j^z S_{j+4}^z \rangle$ .

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# 1 Introduction

The anti-ferromagnetic spin-1/2 Heisenberg  $XXX$  chain [1]

$$H = \sum_{j=1}^L (S_j^x S_{j+1}^x + S_j^y S_{j+1}^y + S_j^z S_{j+1}^z), \quad (1.1)$$

is one of the most fundamental models in the study of quantum magnetism in low dimensions (see, for instance, the book [2]).

The Hamiltonian (1.1) was diagonalized by Bethe in 1931 [3] and the ground state in the thermodynamic limit was investigated by Hulthén in 1938 [4]. The physical content was clarified by Faddeev and Takhtajan [6] using the algebraic Bethe ansatz formulated by Faddeev, Sklyanin and Takhtajan (see the review [5]). Exact integrability of the Heisenberg model is intimately connected with the Yang-Baxter equation [7].

The Heisenberg model also appears as an effective Hamiltonian of more complicated systems such as the one-dimensional Hubbard model at half-filling with a large coupling constant  $U$ .

Since discovering of the model a lot of physical quantities have been calculated exactly which afford significant and reliable data for the experiments as well as numerical simulations. However, most of them are the bulk quantities and we have obtained little results on the correlation functions. Especially significant are the (static) spin-spin pair correlators  $\langle S_j^z S_{j+k}^z \rangle$  (or equivalently  $\langle S_j^+ S_{j+k}^- \rangle / 2$ ;  $S_j^\pm = S_j^x \pm i S_j^y$ ), for which only the first and the second neighbor ( $k = 1, 2$ ) have been known for a long time:

$$\langle S_j^z S_{j+1}^z \rangle = \frac{1}{12} - \frac{1}{3} \ln 2 \simeq -0.14771572685, \quad (1.2)$$

$$\langle S_j^z S_{j+2}^z \rangle = \frac{1}{12} - \frac{4}{3} \ln 2 + \frac{3}{4} \zeta(3) \simeq 0.06067976996. \quad (1.3)$$

Here  $\zeta(s)$  is the Riemann zeta function and  $\langle \dots \rangle$  denotes the ground state expectation value. Note that (1.2) was derived directly from the ground state energy in the thermodynamic limit by Hulthén, while (1.3) was obtained by Takahashi in 1977 via the strong coupling expansion for the ground state energy of the half-filled Hubbard model [8] (see also another derivation by Dittrich and Izozemtsev [9]). Unfortunately we cannot generalize these methods so as to calculate further correlations  $\langle S_j^z S_{j+k}^z \rangle_{k \geq 3}$ .

On the other hand, in 1990s, correlation functions were studied systematically for the Heisenberg  $XXZ$  model. It was mainly developed by Kyoto group (Jimbo, Miwa, Miki, Nakayashiki, [14], see also the book [15]). Utilizing the representation theory of the quantum affine algebra  $U_q(\widehat{sl_2})$ , they first derived multiple integral representation of arbitrary correlators for the massive  $XXZ$  anti-ferromagnet. In 1996 Jimbo and Miwa extended the above results to the  $XXX$  chain and the massless  $XXZ$  chain [17]. The manifest integral formula for special case of correlation functions, Emptiness Formation Probability (EFP) of the  $XXX$  model was earlier obtained by Korepin, Izergin, Essler and Uglov in [16]. Strictly speaking the formulas in the massless region were a conjecture since they were obtained as solutions to the quantum Knizhnik-Zamolodchikov equations (qKZ equations) with a certain assumption. However, the same integral formulas have been reproduced by Kitanine, Maillet and Terras, [20, 21] in the framework of the Quantum Inverse Scattering Method. So we have now at hand the multiple integral representations for any arbitrary correlation functions for the Heisenberg  $XXZ$  chain.

In 2001, Boos and Korepin invented a method to evaluate these multiple integrals for the Heisenberg  $XXX$  chain [22, 23]. They have shown that the integrand can be reduced to a certain canonical form so that the integrations can be performed. They applied the method to the case of the Emptiness Formation Probability. There has been also suggested conjecture that the correlation functions for the  $XXX$  model may be expressed in terms of the Riemann's zeta function at odd arguments with rational coefficients. For instance, the results for  $P(4)$  and  $P(5)$  look as follows:

$$\begin{aligned} P(4) &\equiv \langle (1/2 + S_j^z) (1/2 + S_{j+1}^z) (1/2 + S_{j+2}^z) (1/2 + S_{j+3}^z) \rangle \\ &= \frac{1}{5} - 2 \ln 2 + \frac{173}{60} \zeta(3) - \frac{11}{6} \ln 2 \cdot \zeta(3) - \frac{51}{80} \zeta(3)^2 - \frac{55}{24} \zeta(5) + \frac{85}{24} \ln 2 \cdot \zeta(5) \end{aligned} \tag{1.4}$$

$$\begin{aligned} P(5) &\equiv \langle (1/2 + S_j^z) (1/2 + S_{j+1}^z) (1/2 + S_{j+2}^z) (1/2 + S_{j+3}^z) (1/2 + S_{j+4}^z) \rangle \\ &= \frac{1}{6} - \frac{10}{3} \ln 2 + \frac{281}{24} \zeta(3) - \frac{45}{2} \ln 2 \cdot \zeta(3) - \frac{489}{16} \zeta(3)^2 - \frac{6775}{192} \zeta(5) + \frac{1225}{6} \ln 2 \cdot \zeta(5) \\ &\quad - \frac{425}{64} \zeta(3) \cdot \zeta(5) - \frac{12125}{256} \zeta(5)^2 + \frac{6223}{256} \zeta(7) - \frac{11515}{64} \ln 2 \cdot \zeta(7) + \frac{42777}{512} \zeta(3) \cdot \zeta(7) \end{aligned} \tag{1.5}$$

The last result was obtained in [28].

More recently the method was applied to calculate other correlation functions of the Heisenberg

chain (1.1). Here let us list the results in [29] where all the independent correlation functions on four lattice sites are obtained:

$$\begin{aligned}
\langle S_j^z S_{j+3}^z \rangle &= \frac{1}{12} - 3 \ln 2 + \frac{37}{6} \zeta(3) - \frac{14}{3} \ln 2 \cdot \zeta(3) - \frac{3}{2} \zeta(3)^2 - \frac{125}{24} \zeta(5) + \frac{25}{3} \ln 2 \cdot \zeta(5) \\
&= -0.05024862725 \dots, \\
\langle S_j^x S_{j+1}^x S_{j+2}^z S_{j+3}^z \rangle &= \frac{1}{240} + \frac{1}{12} \ln 2 - \frac{91}{240} \zeta(3) + \frac{1}{6} \ln 2 \cdot \zeta(3) + \frac{3}{80} \zeta(3)^2 + \frac{35}{96} \zeta(5) - \frac{5}{24} \ln 2 \cdot \zeta(5) \\
\langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+3}^z \rangle &= \frac{1}{240} - \frac{1}{6} \ln 2 + \frac{77}{120} \zeta(3) - \frac{5}{12} \ln 2 \cdot \zeta(3) - \frac{3}{20} \zeta(3)^2 - \frac{65}{96} \zeta(5) + \frac{5}{6} \ln 2 \cdot \zeta(5) \\
\langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+3}^x \rangle &= \frac{1}{240} - \frac{1}{4} \ln 2 + \frac{169}{240} \zeta(3) - \frac{5}{12} \ln 2 \cdot \zeta(3) - \frac{3}{20} \zeta(3)^2 - \frac{65}{96} \zeta(5) + \frac{5}{6} \ln 2 \cdot \zeta(5)
\end{aligned} \tag{1.6}$$

Note that the other correlation functions among four lattice sites are expressed as a linear combination of the correlations above. For example we have

$$\begin{aligned}
\langle S_j^z S_{j+1}^z S_{j+2}^z S_{j+3}^z \rangle &= \langle S_j^x S_{j+1}^x S_{j+2}^z S_{j+3}^z \rangle + \langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+3}^z \rangle + \langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+3}^x \rangle \\
&= \frac{1}{80} - \frac{1}{3} \ln 2 + \frac{29}{30} \zeta(3) - \frac{2}{3} \ln 2 \cdot \zeta(3) - \frac{21}{80} \zeta(3)^2 - \frac{95}{96} \zeta(5) + \frac{35}{24} \ln 2 \cdot \zeta(5), \\
P(4) &= \frac{1}{16} + \frac{3}{4} \langle S_j^z S_{j+1}^z \rangle + \frac{1}{2} \langle S_j^z S_{j+2}^z \rangle + \frac{1}{4} \langle S_j^z S_{j+3}^z \rangle + \langle S_j^z S_{j+1}^z S_{j+2}^z S_{j+3}^z \rangle,
\end{aligned} \tag{1.7}$$

etc....

In papers [30, 31, 32] the above results were generalized to the  $XXZ$  Heisenberg chain both in massless and in massive regime.

Generalizing the results (1.6) to the correlation functions on five lattice sites is a natural next problem. In principle, it is possible to calculate them similarly from the multiple integrals as was actually done for  $P(5)$ . It, however, will take a tremendous amount of work to complete them.

It is worthy to note that all the above results have one specific feature, namely, they all are expressed in terms of one-dimensional integrals. In series of papers [25, 26, 27] this phenomenon was explained by means of the above mentioned connection of correlation functions in the inhomogeneous case with the quantum Knizhnik-Zamolodchikov equation, more precisely, by means of a duality between solutions to the qKZ equation on level -4 and 0. It was also argued that this fact has a deep mathematical origin based on theory of deformed hyper-elliptic integrals, symplectic group and special type of cohomologies investigated earlier by Nakayashiki and Smirnov [18, 19].

One of basic outcomes of the above scheme is a general ansatz for correlation functions in which only one transcendental function participates with some rational functions as coefficients. Some special additional properties derived from the qKZ equation allowed to fix  $P(n)$  until  $n = 6$ .

In recent paper [33] the properties of the qKZ equation have been further investigated and the above ansatz has been proved by means of some special recursion relations which linearly connect correlation functions at  $n, n - 1$  and  $n - 2$ . Coefficients in these relations appeared to be related to special transfer matrices over an auxiliary space of “fractional dimension”. The solution to these relations has also been found in the paper [33] which gives an explicit result for rational coefficients before the transcendental functions staying in the ansatz for correlation functions. In spite of the fact that this formula is explicit for any number of sites  $n$ , it is still hard to use it for getting manifest results for concrete  $n$ .

In this paper we generalize the scheme suggested in paper [25] to the case of an arbitrary correlation function. We use a set of relations for correlation functions that follow from the qKZ equation. We optionally call them “first principle” relations. The claim is that these relations together with the above mentioned ansatz completely fix rational functions which participate in this ansatz. We find the whole set of correlation functions in case of  $n = 5$ . After taking homogeneous limit we come to the result in terms of the Riemann zeta function in full agreement with the above conjecture, which now should be considered as proved. The “first principle” relations as well as results for correlation functions should be definitely equivalent to those given by the formula (3.22) in the paper [33]. But it is still to show such an equivalence explicitly.

## 2 Quantum Knizhnik-Zamolodchikov equations and fundamental relations for correlation functions

In this section we would like to describe an algebraic scheme which allows to calculate any correlation function for the  $XXX$  model. In fact we generalize here the method applied in paper [24] to the particular case of the EFP.

Let  $R(l)$  be the  $R$ -matrix for the  $XXX$  model

$$R(l) = \frac{R_0(l)}{l + \pi i} \begin{pmatrix} l + \pi i & 0 & 0 & 0 \\ 0 & l & \pi i & 0 \\ 0 & \pi i & l & 0 \\ 0 & 0 & 0 & l + \pi i \end{pmatrix}, \quad (2.1)$$

where

$$R_0(l) = -\frac{\Gamma\left(\frac{l}{2\pi i}\right)\Gamma\left(\frac{1}{2} - \frac{l}{2\pi i}\right)}{\Gamma\left(-\frac{l}{2\pi i}\right)\Gamma\left(\frac{1}{2} + \frac{l}{2\pi i}\right)}.$$

Sometimes we shall also use notation

$$\overline{R}(z) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{z}{z+i} & \frac{i}{z+i} & 0 \\ 0 & \frac{i}{z+i} & \frac{z}{z+i} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad (2.2)$$

where it is implied that  $l = \pi z$ .

The corner stone of our scheme is the quantum Knizhnik-Zamolodchikov equation [qKZ] [10, 13, 11, 12]. It is a linear system of equations for a vector function  $g_n$ . In particular, the qKZ equation on level  $-4$  looks as follows

$$\begin{aligned} g_n(l_1, \dots, l_{j+1}, l_j, \dots, l_{2n})_{\epsilon_1, \dots, \epsilon'_{j+1}, \epsilon'_j, \dots, \epsilon_{2n}} &= \\ &= R(l_j - l_{j+1})_{\epsilon'_j, \epsilon'_{j+1}}^{\epsilon_j, \epsilon_{j+1}} g_n(l_1, \dots, l_j, l_{j+1}, \dots, l_{2n})_{\epsilon_1, \dots, \epsilon_j, \epsilon_{j+1}, \dots, \epsilon_{2n}} \end{aligned} \quad (2.3)$$

$$g_n(l_1, \dots, l_{2n-1}, l_{2n} + 2\pi i)_{\epsilon_1, \dots, \epsilon_{2n-1}, \epsilon_{2n}} = g_n(l_{2n}, l_1, \dots, l_{2n-1})_{\epsilon_{2n}, \epsilon_1, \dots, \epsilon_{2n-1}} \quad (2.4)$$

It was suggested by the Kyoto school [14] that the correlation functions are connected with solutions to the qKZ equation. In particular, for the massless case Jimbo and Miwa [17] found that one should take a special solution of the qKZ on level -4 which satisfies one additional constraint

$$\begin{aligned} g_n(l_1, \dots, l_{j-1}, l_j, l_j - \pi i, l_{j+2}, \dots, l_{2n})_{\epsilon_1, \dots, \epsilon_{j-1}, \epsilon_j, \epsilon_{j+1}, \epsilon_{j+2}, \dots, \epsilon_{2n}} &= \\ &= \epsilon_j \delta_{\epsilon_j, -\epsilon_{j+1}} g_{n-1}(l_1, \dots, l_{j-1}, l_{j+2}, \dots, l_{2n})_{\epsilon_1, \dots, \epsilon_{j-1}, \epsilon_{j+2}, \dots, \epsilon_{2n}} \end{aligned} \quad (2.5)$$

Solutions to the above equations (2.4,2.5) are meromorphic functions with possible singularities at the points

$$\Im(l_j - l_k) = \pi l, \quad l \in \mathbf{Z} \setminus 0$$

but regular at  $\Im(l_j - l_k) = \pm\pi$ .

Due to Jimbo and Miwa, any correlation function

$$P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n} \equiv \left\langle E_{\epsilon_1}^{\epsilon'_1} \dots E_{\epsilon_n}^{\epsilon'_n} \right\rangle$$

is given by the following formula

$$P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1, \dots, z_n) = g_n(l_1, \dots, l_n, l_n + \pi i, \dots, l_1 + \pi i)_{\epsilon_1, \dots, \epsilon_n, -\epsilon'_n, \dots, -\epsilon'_1} \quad (2.6)$$

where as above  $l_j = \pi z_j$ .

Let us list the “first principle” relations for the correlation functions which follow from the qKZ equation:

- Translational invariance

$$P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1 + x, \dots, z_n + x) = P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1, \dots, z_n) \quad (2.7)$$

- Transposition, Negating and Reverse order relations

$$\begin{aligned} P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1, \dots, z_n) &= P_{\epsilon'_1, \dots, \epsilon'_n}^{\epsilon_1, \dots, \epsilon_n}(-z_1, \dots, -z_n) = \\ P_{-\epsilon_1, \dots, -\epsilon_n}^{-\epsilon'_1, \dots, -\epsilon'_n}(z_1, \dots, z_n) &= P_{\epsilon_n, \dots, \epsilon_1}^{\epsilon'_n, \dots, \epsilon'_1}(-z_n, \dots, -z_1) \end{aligned} \quad (2.8)$$

- Intertwining relation

$$\begin{aligned} \overline{R}_{\tilde{\epsilon}'_j \tilde{\epsilon}'_{j+1}}^{\epsilon'_j \epsilon'_{j+1}}(z_j - z_{j+1}) P_{\dots \epsilon_{j+1}, \epsilon_j \dots}^{\dots \tilde{\epsilon}'_{j+1}, \tilde{\epsilon}'_j \dots}(\dots z_{j+1}, z_j \dots) = \\ P_{\dots \tilde{\epsilon}_j, \tilde{\epsilon}_{j+1} \dots}^{\dots \epsilon'_j, \epsilon'_{j+1} \dots}(\dots z_j, z_{j+1} \dots) \overline{R}_{\epsilon_j \epsilon_{j+1}}^{\tilde{\epsilon}_j \tilde{\epsilon}_{j+1}}(z_j - z_{j+1}) \end{aligned} \quad (2.9)$$

- First recurrent relation

$$\begin{aligned} P_{\epsilon_1, \epsilon_2, \dots, \epsilon_n}^{\epsilon'_1, \epsilon'_2, \dots, \epsilon'_n}(z+i, z, z_3, \dots, z_n) &= -\delta_{\epsilon_1, -\epsilon_2} \epsilon'_1 \epsilon_2 P_{-\epsilon'_1, \epsilon_3, \dots, \epsilon_n}^{\epsilon'_2, \epsilon'_3, \dots, \epsilon'_n}(z, z_3, \dots, z_n) \\ P_{\epsilon_1, \epsilon_2, \dots, \epsilon_n}^{\epsilon'_1, \epsilon'_2, \dots, \epsilon'_n}(z-i, z, z_3, \dots, z_n) &= -\delta_{\epsilon'_1, -\epsilon'_2} \epsilon'_1 \epsilon'_2 P_{\epsilon_2, \epsilon_3, \dots, \epsilon_n}^{-\epsilon_1, \epsilon'_3, \dots, \epsilon'_n}(z, z_3, \dots, z_n) \end{aligned} \quad (2.10)$$

- Second recurrent relation

$$\lim_{z_j \rightarrow \infty} P_{\epsilon_1, \dots, \epsilon_j, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_j, \dots, \epsilon'_n}(z_1, \dots, z_j, \dots, z_n) = \delta_{\epsilon_j, \epsilon'_j} \frac{1}{2} P_{\epsilon_1, \dots, \hat{\epsilon}_j, \dots, \epsilon_n}^{\epsilon'_1, \dots, \hat{\epsilon}'_j, \dots, \epsilon'_n}(z_1, \dots, \hat{z}_j, \dots, z_n) \quad (2.11)$$

- Identity relation

$$\begin{aligned} \sum_{\substack{\epsilon_1, \dots, \epsilon_n \\ \sum_i \epsilon'_i = \sum_i \epsilon_i}} P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1, \dots, z_n) &= \sum_{\substack{\epsilon'_1, \dots, \epsilon'_n \\ \sum_i \epsilon'_i = \sum_i \epsilon_i}} P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1, \dots, z_n) = \\ &= P_{+, \dots, +}^{+, \dots, +}(z_1, \dots, z_n) = P_{-, \dots, -}^{-, \dots, -}(z_1, \dots, z_n) \end{aligned} \quad (2.12)$$

- Reduction relation

$$\begin{aligned} P_{+, \epsilon_2, \dots, \epsilon_n}^{+, \epsilon'_2, \dots, \epsilon'_n}(z_1, z_2, \dots, z_n) + P_{-, \epsilon_2, \dots, \epsilon_n}^{-, \epsilon'_2, \dots, \epsilon'_n}(z_1, z_2, \dots, z_n) &= P_{\epsilon_2, \dots, \epsilon_n}^{\epsilon'_2, \dots, \epsilon'_n}(z_2, \dots, z_n) \\ P_{\epsilon_1, \dots, \epsilon_{n-1}, +}^{\epsilon'_1, \dots, \epsilon'_{n-1}, +}(z_1, z_2, \dots, z_n) + P_{\epsilon_1, \dots, \epsilon_{n-1}, -}^{\epsilon'_1, \dots, \epsilon'_{n-1}, -}(z_1, z_2, \dots, z_n) &= P_{\epsilon_1, \dots, \epsilon_{n-1}}^{\epsilon'_1, \dots, \epsilon'_{n-1}}(z_1, \dots, z_{n-1}) \end{aligned} \quad (2.13)$$

The relations (2.7), (2.8) and (2.13) were discussed by Jimbo and Miwa in the paper [17]. The intertwining relation (2.9) is the direct corollary of the qKZ equations. The proof of the first and the second recurrent relations (2.10,2.11) can be performed in the same lines as it was described in the paper [24]. In particular, the relation (2.11) was completely proved in the paper [33]. The identity relation (2.12) is a direct consequence of the  $sl_2$  symmetry of the solution to the qKZ.

In papers [25, 26], a general ansatz for the correlation functions was suggested. Then it was proved<sup>1</sup> in [33]

$$P_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(z_1, \dots, z_n) = \sum_{m=0}^{\lfloor \frac{n}{2} \rfloor} \sum_{1 \leq k_1 \neq \dots \neq k_{2m} \leq n} A_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(k_1, \dots, k_{2m} | z_1, \dots, z_n) G(z_{k_1} - z_{k_2}) \cdots G(z_{k_{2m-1}} - z_{k_{2m}}) \quad (2.14)$$

where

$$G(x) = 2 \sum_{k=1}^{\infty} (-1)^k \cdot k \cdot \frac{x^2 + 1}{x^2 + k^2} \quad (2.15)$$

and  $A_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(k_1, \dots, k_{2m} | z_1, \dots, z_n)$  are rational functions with known denominator and polynomial numerator of a known degree in variables  $z_1, \dots, z_n$ , namely,

$$A_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(k_1, \dots, k_{2m} | z_1, \dots, z_n) = \frac{Q_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(k_1, \dots, k_{2m} | z_1, \dots, z_n)}{\prod'_{i < j} (z_i - z_j)} \quad (2.16)$$

where  $\prod'_{i < j} (z_i - z_j)$  means the product in which the differences with  $\{i, j\} = \{k_{2l}, k_{2l+1}\}$  and  $i, j \notin K$  should be excluded,  $K$  is the set  $k_1, \dots, k_{2m}$ .  $Q_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(k_1, \dots, k_{2m} | z_1, \dots, z_n)$  are polynomials of the variables  $z$  which have the same degree in each variable as in the denominator and also the same degree of homogeneity. Some useful properties of the function  $G(x)$  are listed below [24],

$$G(-x) = G(x), \quad (2.17)$$

$$G(\pm\infty) = -\frac{1}{2}, \quad (2.18)$$

$$G(x - i) = \alpha(x) + \gamma(x)G(x), \quad (2.19)$$

where

$$\alpha(x) = -\frac{x - 2i}{x - i}, \quad \gamma(x) = -\frac{x(x - 2i)}{x^2 + 1}. \quad (2.20)$$

Moreover  $G(x)$  is a generating function of the alternating zeta functions

$$G(x) = -2(1 + x^2) \sum_{k=0}^{\infty} (-1)^k x^{2k} \zeta_a(2k + 1), \quad (2.21)$$

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<sup>1</sup>The function  $\omega(l)$  used in [33] is related to the function  $G$  by  $\omega(l) = G(il) + 1/2$

with

$$\zeta_a(s) \equiv \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n^s} = (1 - 2^{1-s})\zeta(s). \quad (2.22)$$

Conjectural formula for the very first term when  $m = 0$  looks as follows

$$A_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(0|z_1, \dots, z_n) = Q_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(0|z_1, \dots, z_n) = \frac{1}{(n+1)\binom{n}{r}} \quad (2.23)$$

where  $r$  is the number of down (or up) spins in the set  $\epsilon_1, \dots, \epsilon_n$  (or  $\epsilon'_1, \dots, \epsilon'_n$ ).

Our claim is that the “first principle” relations (2.7)-(2.13) together with the ansatz (2.14) and conjecture (2.23) completely defines all unknown coefficients in polynomials  $Q_{\epsilon_1, \dots, \epsilon_n}^{\epsilon'_1, \dots, \epsilon'_n}(k_1, \dots, k_{2m}|z_1, \dots, z_n)$  and hence the whole set of correlation functions.

Here let us list the explicit form of the inhomogeneous correlation functions for  $n = 1$  and  $n = 2$ .

$$\begin{aligned} P_+^+(z_1) &= P_-^-(z_1) = \frac{1}{2}, \\ P_{++}^{++}(z_1, z_2) &= P_{--}^{--}(z_1, z_2) = \frac{1}{3} + \frac{1}{6}G(z_1 - z_2), \\ P_{+-}^{+-}(z_1, z_2) &= P_{-+}^{-+}(z_1, z_2) = \frac{1}{6} - \frac{1}{6}G(z_1 - z_2), \\ P_{+-}^{-+}(z_1, z_2) &= P_{-+}^{+-}(z_1, z_2) = \frac{1}{6} + \frac{1}{3}G(z_1 - z_2). \end{aligned} \quad (2.24)$$

One can easily confirm these correlation functions satisfy the first principle relations.

In the subsequent sections, we denote the inhomogeneous EFP as

$$P_n(z_1, z_2, \dots, z_n) \equiv P_{++ \dots +}^{++ \dots +}(z_1, z_2, \dots, z_n) \left( = P_{-- \dots -}^{-- \dots -}(z_1, z_2, \dots, z_n) \right). \quad (2.25)$$

Especially we have  $P_2(z_1, z_2) = P_{++}^{++}(z_1, z_2)$ . The explicit form of the above inhomogeneous EFPs up to  $n = 6$  was obtained in [24].

### 3 Inhomogeneous correlation functions for $n = 3$

Here we consider the inhomogeneous correlation functions for  $n = 3$  and the corresponding first principle relations. This case was considered as an example in the paper [33]. But we reproduce

$P_{\epsilon_1 \epsilon_2 \epsilon_3}^{\epsilon'_1 \epsilon'_2 \epsilon'_3}$	— + +	+ — +	+ + —
— + +	$\frac{1}{12} + \frac{1}{12}(G_{23} - G_{12} - G_{13}) + \frac{1}{12}G_{123}$	$\frac{1}{12} + \frac{1}{6}G_{12} + \frac{i}{12}G_{123}^A$	$\frac{1}{12} + \frac{1}{6}G_{13} - \frac{1}{6}G_{123} - \frac{i}{12}G_{123}^A$
+ — +	$\frac{1}{12} + \frac{1}{6}G_{12} - \frac{i}{12}G_{123}^A$	$\frac{1}{12} + \frac{1}{12}(G_{13} - G_{23} - G_{12}) - \frac{1}{12}G_{123}$	$\frac{1}{12} + \frac{1}{6}G_{23} + \frac{i}{12}G_{123}^A$
+ + —	$\frac{1}{12} + \frac{1}{6}G_{13} - \frac{1}{6}G_{123} + \frac{i}{12}G_{123}^A$	$\frac{1}{12} + \frac{1}{6}G_{23} - \frac{i}{12}G_{123}^A$	$\frac{1}{12} + \frac{1}{12}(G_{12} - G_{13} - G_{23}) + \frac{1}{12}G_{123}$

Table 1: Inhomogeneous correlation functions for  $n = 3$  in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 = 1$

it here in order to make our approach more illustrative. To describe the explicit expressions of  $P_{\epsilon_1 \epsilon_2 \epsilon_3}^{\epsilon'_1 \epsilon'_2 \epsilon'_3}(z_1, z_2, z_3)$ , we introduce some notations.

$$G_{ij} \equiv G(z_{ij}), \quad G_{123} \equiv \frac{G_{12}}{z_{13}z_{23}} + \frac{G_{13}}{z_{12}z_{32}} + \frac{G_{23}}{z_{21}z_{31}},$$

$$G_{123}^A = \left( \frac{1}{z_{13}} - \frac{1}{z_{23}} \right) G_{12} + \left( \frac{1}{z_{23}} - \frac{1}{z_{12}} \right) G_{13} + \left( \frac{1}{z_{31}} - \frac{1}{z_{21}} \right) G_{23}, \quad (3.1)$$

where we have also introduced an abbreviation  $z_{ij} = z_i - z_j$ . With these notations, the EFP  $P_3(z_1, z_2, z_3) \equiv P_{+++}^{+++}(z_1, z_2, z_3)$  is given by [24]

$$P_3(z_1, z_2, z_3) = \frac{1}{4} + \frac{1}{12}(G_{12} + G_{13} + G_{23}) - \frac{1}{12}G_{123}. \quad (3.2)$$

Next we consider the inhomogeneous correlation functions  $P_{\epsilon_1 \epsilon_2 \epsilon_3}^{\epsilon'_1 \epsilon'_2 \epsilon'_3}(z_1, z_2, z_3)$  in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 = 1$ . For convenience, we first exhibit the final results in Table 1. There each row corresponds to the subscript  $\epsilon_1, \epsilon_2, \epsilon_3$ , while each column corresponds to the superscript  $\epsilon'_1, \epsilon'_2, \epsilon'_3$ .

Below we show the strategy to calculate these correlation functions from our first principle relations. Assume we have already obtained the EFP,  $P_3(z_1, z_2, z_3)$  and all the correlations for  $n = 2$ . As is shown in Table 1, there are 9 correlation functions in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 = 1$ .

However, we find, they are actually connected through the intertwining relation (2.9) as follows,

$$\begin{aligned}
P_{+-+}^{--+}(z_1, z_2, z_3) &= \frac{z_{23} + i}{z_{23}} P_{++-}^{--+}(z_1, z_3, z_2) - \frac{i}{z_{23}} P_{++-}^{--+}(z_1, z_2, z_3), \\
P_{++-}^{+-+}(z_1, z_2, z_3) &= \frac{z_{21} + i}{z_{21}} P_{++-}^{--+}(z_2, z_1, z_3) - \frac{i}{z_{21}} P_{++-}^{--+}(z_1, z_2, z_3), \\
P_{-++}^{+-+}(z_1, z_2, z_3) &= P_{-++}^{--+}(z_2, z_1, z_3) + \frac{i}{z_{21}} \{ P_{-++}^{--+}(z_2, z_1, z_3) - P_{-++}^{--+}(z_1, z_2, z_3) \}, \\
P_{-++}^{++-}(z_1, z_2, z_3) &= \frac{z_{32} + i}{z_{32}} P_{-++}^{++-}(z_1, z_3, z_2) - \frac{i}{z_{32}} P_{-++}^{++-}(z_1, z_2, z_3), \\
P_{-++}^{+-+}(z_1, z_2, z_3) &= P_{-++}^{--+}(z_2, z_1, z_3) + \frac{i}{z_{21}} \{ P_{-++}^{--+}(z_2, z_1, z_3) - P_{-++}^{--+}(z_1, z_2, z_3) \}, \\
P_{-++}^{++-}(z_1, z_2, z_3) &= \frac{z_{12} + i}{z_{12}} P_{-++}^{++-}(z_2, z_1, z_3) - \frac{i}{z_{12}} P_{-++}^{++-}(z_1, z_2, z_3), \\
P_{++-}^{++-}(z_1, z_2, z_3) &= P_{++-}^{--+}(z_1, z_3, z_2) + \frac{i}{z_{32}} \{ P_{++-}^{--+}(z_1, z_3, z_2) - P_{++-}^{--+}(z_1, z_2, z_3) \}. \tag{3.3}
\end{aligned}$$

These relations imply we need to calculate only two correlations, for example,  $P_{-++}^{--+}(z_1, z_2, z_3)$  and  $P_{++-}^{--+}(z_1, z_2, z_3)$ . The remaining correlation functions are obtained immediately from (3.3). Moreover the diagonal one,  $P_{-++}^{--+}(z_1, z_2, z_3)$  can be derived easily from the reduction relation (2.13),

$$\begin{aligned}
P_{-++}^{--+}(z_1, z_2, z_3) &= P_2(z_2, z_3) - P_3(z_1, z_2, z_3) \\
&= \frac{1}{3} + \frac{1}{6} G_{23} - \left( \frac{1}{4} + \frac{1}{12} (G_{12} + G_{13} + G_{23}) - \frac{1}{12} G_{123} \right) \\
&= \frac{1}{12} + \frac{1}{12} (G_{23} - G_{12} - G_{13}) + \frac{1}{12} G_{123}. \tag{3.4}
\end{aligned}$$

Therefore our task reduces to calculate one off-diagonal correlation  $P_{++-}^{--+}(z_1, z_2, z_3)$ . Note that there are different ways to connect the correlations from (3.3). They, however, are consistent with each other due to the Yang-Baxter relation for  $\bar{R}_{ij}(z)$ .

Look at other first principle relations and apply them to  $P_{++-}^{--+}(z_1, z_2, z_3)$ . First, we find the reverse order relation gives

$$P_{++-}^{--+}(z_1, z_2, z_3) = P_{++-}^{--+}(z_3, z_2, z_1), \tag{3.5}$$

which means  $P_{++-}^{--+}(z_1, z_2, z_3)$  is symmetric with respect to the exchange  $z_1 \leftrightarrow z_3$ . Next, the transposition relation  $P_{-++}^{++-}(z_1, z_2, z_3) = P_{-++}^{++-}(-z_1, -z_2, -z_3)$  together with (3.3), (3.4) and (3.5)

gives an equation

$$\begin{aligned}
& P_{++-}^{---}(-z_1, -z_2, -z_3) \\
&= \frac{z_{12} + i}{z_{12}} \frac{z_{32} + i}{z_{32}} P_{++-}^{---}(z_1, z_2, z_3) - \frac{i}{z_{12}} \frac{z_{31} + i}{z_{31}} P_{++-}^{---}(z_2, z_1, z_3) - \frac{i}{z_{32}} \frac{z_{13} + i}{z_{13}} P_{++-}^{---}(z_1, z_3, z_2) \\
&\quad - \frac{1}{z_{12} z_{32}} P_2(z_1, z_3) + \frac{i}{z_{31}} \frac{z_{32} + i}{z_{32}} P_2(z_1, z_2) + \frac{i}{z_{13}} \frac{z_{12} + i}{z_{12}} P_2(z_2, z_3).
\end{aligned} \tag{3.6}$$

The recurrent relations for  $P_{++-}^{---}(z_1, z_2, z_3)$  are given by

$$P_{++-}^{---}(z + i, z, z_3) = 0, \quad P_{++-}^{---}(z - i, z, z_3) = -P_{+-}^{--}(z, z_3), \tag{3.7}$$

$$\begin{aligned}
\lim_{z_1 \rightarrow \infty} P_{++-}^{---}(z_1, z_2, z_3) &= \lim_{z_3 \rightarrow \infty} P_{++-}^{---}(z_1, z_2, z_3) = 0, \\
\lim_{z_2 \rightarrow \infty} P_{++-}^{---}(z_1, z_2, z_3) &= \frac{1}{2} P_{+-}^{--}(z_1, z_3).
\end{aligned} \tag{3.8}$$

The identity relation gives a relation,

$$\frac{z_{32} + i}{z_{32}} P_{++-}^{---}(z_1, z_2, z_3) + \frac{z_{23} + i}{z_{23}} P_{++-}^{---}(z_1, z_3, z_2) = 2P_3(z_1, z_2, z_3) - P_2(z_2, z_3). \tag{3.9}$$

where we have used the relations (3.3) and (3.4). Equations (3.5)–(3.9) with the translational invariance

$$P_{++-}^{---}(z_1 + x, z_2 + x, z_3 + x) = P_{++-}^{---}(z_1, z_2, z_3). \tag{3.10}$$

are our first principle relations reduced to  $P_{++-}^{---}(z_1, z_2, z_3)$ .

According to the ansatz (2.14) and (2.23), let us assume  $P_{++-}^{---}(z_1, z_2, z_3)$  in the form

$$\begin{aligned}
P_{++-}^{---}(z_1, z_2, z_3) &= \frac{1}{12} + \frac{c_{11} + c_{12}z_{13} + c_{13}z_{23} + c_{14}z_{13}z_{23}}{z_{13}z_{23}} G_{12} \\
&\quad + \frac{c_{21} + c_{22}z_{12} + c_{23}z_{32} + c_{24}z_{12}z_{32}}{z_{12}z_{32}} G_{13} \\
&\quad + \frac{c_{31} + c_{32}z_{21} + c_{33}z_{31} + c_{34}z_{21}z_{31}}{z_{21}z_{31}} G_{23},
\end{aligned} \tag{3.11}$$

where  $c_{ij}$  are the coefficients which will be determined by the relations (3.5)–(3.9). Note that from the symmetry relation (3.5), we at once have several relations

$$c_{31} = c_{11}, \quad c_{32} = c_{13}, \quad c_{33} = c_{12}, \quad c_{34} = c_{14}, \quad c_{23} = c_{22}. \tag{3.12}$$

Relations (3.6)–(3.9) generate equations for the coefficients  $c_{ij}$ . Although they are an overdetermined system of the equations for  $c_{ij}$ , one can figure out it has a unique solution

$$\begin{aligned} c_{11} = c_{21} = c_{31} &= -\frac{1}{6}, \quad c_{12} = -c_{13} = c_{22} = c_{23} = -c_{32} = c_{33} = -\frac{i}{12}, \\ c_{14} = c_{34} &= 0, \quad c_{24} = \frac{1}{6}. \end{aligned} \tag{3.13}$$

Namely the correlation function  $P_{++-}^{--+}(z_1, z_2, z_3)$  is determined as

$$\begin{aligned} P_{++-}^{--+}(z_1, z_2, z_3) &= \frac{1}{12} + \frac{1}{6} \left\{ -\frac{1}{z_{13}z_{23}} + \frac{i}{2} \left( -\frac{1}{z_{23}} + \frac{1}{z_{13}} \right) \right\} G_{12} \\ &\quad + \frac{1}{6} \left\{ 1 - \frac{1}{z_{12}z_{32}} - \frac{i}{2} \left( \frac{1}{z_{32}} + \frac{1}{z_{12}} \right) \right\} G_{13} \\ &\quad + \frac{1}{6} \left\{ -\frac{1}{z_{21}z_{31}} + \frac{i}{2} \left( \frac{1}{z_{31}} - \frac{1}{z_{21}} \right) \right\} G_{23} \\ &= \frac{1}{12} + \frac{1}{6}G_{13} - \frac{1}{6}G_{123} + \frac{i}{12}G_{123}^A. \end{aligned} \tag{3.14}$$

Other correlation functions are obtained from the intertwining relations (3.3). In this way we could calculate all the inhomogeneous correlation functions as shown in Table 1. The correlation functions in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 = -1$  are simply obtained by the negation relations,

$$P_{\epsilon_1\epsilon_2\epsilon_3}^{\epsilon'_1\epsilon'_2\epsilon'_3}(z_1, z_2, z_3) = P_{-\epsilon_1,-\epsilon_2,-\epsilon_3}^{-\epsilon'_1,-\epsilon'_2,-\epsilon'_3}(z_1, z_2, z_3). \tag{3.15}$$

Now let us think of an inhomogeneous generalization of the next-nearest neighbor correlation function,

$$\begin{aligned} \langle S_1^z S_3^z \rangle(z_1, z_2, z_3) &= \frac{1}{2} \left\{ P_3(z_1, z_2, z_3) - P_{++-}^{--+}(z_1, z_2, z_3) + P_{+-+}^{+-+}(z_1, z_2, z_3) - P_{-++}^{-++}(z_1, z_2, z_3) \right\} \\ &= \frac{1}{4} - P_2(z_1, z_2) - P_2(z_2, z_3) + 2P_3(z_1, z_2, z_3) = \frac{1}{12} + \frac{1}{6}G_{13} - \frac{1}{6}G_{123} \\ &= \frac{1}{12} + \frac{1}{6}G(z_{13}) - \frac{1}{6} \left\{ \frac{G(z_{12})}{z_{13}z_{23}} + \frac{G(z_{13})}{z_{12}z_{32}} + \frac{G(z_{23})}{z_{21}z_{31}} \right\}. \end{aligned} \tag{3.16}$$

Expanding  $G$ -function as

$$G(z) = -2(1 + z^2) \left\{ \zeta_a(1) - z^2 \zeta_a(3) + \dots \right\}, \tag{3.17}$$

and substituting it into (3.16), one can take the homogeneous limit  $z_1, z_2, z_3 \rightarrow 0$ ,

$$\begin{aligned} \lim_{z_i \rightarrow 0} \langle S_1^z S_3^z \rangle(z_1, z_2, z_3) &= \frac{1}{12} - \frac{1}{3} \zeta_a(1) - (\zeta_a(1) - \zeta_a(3)) \\ &= \frac{1}{12} - \frac{4}{3} \ln 2 + \frac{3}{4} \zeta(3), \end{aligned} \quad (3.18)$$

which reproduces the expression (1.3)

## 4 Inhomogeneous correlation functions for $n = 4$

Next we consider the inhomogeneous correlation functions for  $n = 4$ ,  $P_{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4}^{\epsilon'_1 \epsilon'_2 \epsilon'_3 \epsilon'_4}(z_1, z_2, z_3, z_4)$ .

The EFP  $P_4(z_1, z_2, z_3, z_4) = P_{++++}^{++++}(z_1, z_2, z_3, z_4)$  was calculated in [24],

$$\begin{aligned} P_4(z_1, z_2, z_3, z_4) &= \frac{1}{5} + A_{4,1}^{(0)}(z_1, z_2, z_3, z_4)G(z_{12}) + A_{4,1}^{(0)}(z_1, z_3, z_2, z_4)G(z_{13}) \\ &\quad + A_{4,1}^{(0)}(z_1, z_4, z_3, z_2)G(z_{14}) + A_{4,1}^{(0)}(z_3, z_2, z_1, z_4)G(z_{23}) \\ &\quad + A_{4,1}^{(0)}(z_4, z_2, z_3, z_1)G(z_{24}) + A_{4,1}^{(0)}(z_4, z_3, z_2, z_1)G(z_{34}) \\ &\quad + A_{4,2}^{(0)}(z_1, z_2, z_3, z_4)G(z_{12})G(z_{34}) + A_{4,2}^{(0)}(z_1, z_3, z_2, z_4)G(z_{13})G(z_{24}) \\ &\quad + A_{4,2}^{(0)}(z_1, z_4, z_3, z_2)G(z_{14})G(z_{23}), \end{aligned}$$

with

$$\begin{aligned} A_{4,i}^{(0)}(z_1, z_2, z_3, z_4) &= \frac{Q_{4,i}^{(0)}(z_1, z_2, z_3, z_4)}{z_{13}z_{14}z_{23}z_{24}}, \quad (i = 1, 2), \\ Q_{4,2}^{(0)}(z_1, z_2, z_3, z_4) &= \frac{1}{36} \left\{ (z_{13}z_{24} - 1)(z_{14}z_{23} - 1) + \frac{2}{5} \left( z_{12}^2 + \frac{3}{2} \right) \left( z_{34}^2 + \frac{3}{2} \right) + \frac{3}{2} \right\}, \\ Q_{4,1}^{(0)}(z_1, z_2, z_3, z_4) &= 2Q_{4,2}^{(0)}(z_1, z_2, z_3, z_4) - \frac{1}{60}(z_{12}^2 + 4)(z_{34}^2 + 1). \end{aligned} \quad (4.1)$$

Let us think of the correlation functions in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 + \epsilon_4 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 + \epsilon'_4 = 2$ .

There are 16 correlations in this sector. They, however, are connected via the intertwining relations similarly in the case of  $n = 3$ . In fact, if one can calculate a single correlation, for example,  $P_{+-+-}^{----}(z_1, z_2, z_3, z_4)$ , other correlation functions are obtained with the help of the intertwining relations and the reduction relation

$$P_{-+++}^{----}(z_1, z_2, z_3, z_4) = P_3(z_2, z_3, z_4) - P_4(z_1, z_2, z_3, z_4). \quad (4.2)$$

Therefore we can invoke the same strategy as was done for the calculation of  $P_{++-}^{--+}(z_1, z_2, z_3)$  in the previous section. Namely, we reduce the first principle relations to those for a single functions  $P_{+++-}^{---}(z_1, z_2, z_3, z_4)$  and solve them. Below we list the equations corresponding to (3.5)–(3.9) for  $n = 3$ .

- symmetry relations

$$P_{+++-}^{---}(z_1, z_2, z_3, z_4) = P_{+++-}^{---}(z_1, z_3, z_2, z_4) = P_{+++-}^{---}(z_4, z_3, z_2, z_1) \quad (4.3)$$

- transposition relation

$$\begin{aligned} & P_{+++-}^{---}(-z_1, -z_2, -z_3, -z_4) \\ &= \frac{(z_{12} + i)(z_{13} + i)(z_{42} + i)(z_{43} + i)}{z_{12}z_{13}z_{42}z_{43}} P_{+++-}^{---}(z_1, z_2, z_3, z_4) \\ & - \frac{i(z_{12} + i)(z_{14} + i)(z_{32} + i)}{z_{12}z_{14}z_{32}z_{43}} P_{+++-}^{---}(z_1, z_2, z_4, z_3) \\ & - \frac{i(z_{13} + i)(z_{14} + i)(z_{23} + i)}{z_{13}z_{14}z_{23}z_{42}} P_{+++-}^{---}(z_1, z_3, z_4, z_2) \\ & - \frac{i(z_{43} + i)(z_{41} + i)(z_{23} + i)}{z_{43}z_{41}z_{23}z_{12}} P_{+++-}^{---}(z_2, z_1, z_3, z_4) \\ & - \frac{i(z_{42} + i)(z_{41} + i)(z_{32} + i)}{z_{42}z_{41}z_{32}z_{13}} P_{+++-}^{---}(z_3, z_1, z_2, z_4) \\ & + \frac{1 + i(z_{12} + z_{43}) + z_{12}z_{34} + z_{13}z_{24}}{z_{12}z_{13}z_{24}z_{34}} P_{+++-}^{---}(z_3, z_1, z_2, z_4) \\ & + \frac{i(z_{42} + i)(z_{43} + i)}{z_{42}z_{43}z_{41}} P_3(z_1, z_2, z_3) + \frac{i(z_{12} + i)(z_{13} + i)}{z_{12}z_{13}z_{14}} P_3(z_2, z_3, z_4) + \frac{z_{32} + i}{z_{32}z_{43}z_{31}} P_3(z_1, z_2, z_4) \\ & + \frac{z_{23} + i}{z_{23}z_{42}z_{21}} P_3(z_2, z_3, z_4) \end{aligned} \quad (4.4)$$

- first recurrent relations

$$\begin{aligned} & \lim_{z_1 \rightarrow \infty} P_{+++-}^{---}(z_1, z_2, z_3, z_4) = 0, \\ & \lim_{z_2 \rightarrow \infty} P_{+++-}^{---}(z_1, z_2, z_3, z_4) = \frac{1}{2} P_{++-}^{--+}(z_1, z_3, z_4) \end{aligned} \quad (4.5)$$

- second recurrent relation

$$\begin{aligned} & P_{+++-}^{---}(z + i, z, z_3, z_4) = 0, \\ & P_{+++-}^{---}(z - i, z, z_3, z_4) = -P_{++-}^{--+}(z, z_3, z_4) \end{aligned} \quad (4.6)$$

- identity relation

$$\begin{aligned} & \frac{z_{42} + i}{z_{42}} \frac{z_{43} + i}{z_{43}} P_{+++-}^{----}(z_1, z_2, z_3, z_4) + \frac{z_{23} + i}{z_{23}} \frac{z_{24} + i}{z_{24}} P_{+++-}^{----}(z_1, z_3, z_4, z_2) \\ & + \frac{z_{32} + i}{z_{32}} \frac{z_{34} + i}{z_{34}} P_{+++-}^{----}(z_1, z_2, z_4, z_3) = 2P_4(z_1, z_2, z_3, z_4) - P_3(z_2, z_3, z_4) \end{aligned} \quad (4.7)$$

- translational invariance

$$P_{+++-}^{----}(z_1 + x, z_2 + x, z_3 + x, z_4 + x) = P_{+++-}^{----}(z_1, z_2, z_3, z_4) \quad (4.8)$$

According to the ansatz (2.14), (2.23),  $P_{+++-}^{----}(z_1, z_2, z_3, z_4)$  may be written in the form

$$\begin{aligned} P_{+++-}^{----}(z_1, z_2, z_3, z_4) &= \frac{1}{20} + A_{4,1}^{(1)}(z_1, z_2, z_3, z_4)G(z_{12}) + A_{4,1}^{(1)}(z_1, z_3, z_2, z_4)G(z_{13}) \\ &+ A_{4,2}^{(1)}(z_1, z_4, z_2, z_3)G(z_{14}) + A_{4,3}^{(1)}(z_3, z_2, z_1, z_4)G(z_{23}) \\ &+ A_{4,1}^{(1)}(z_4, z_2, z_3, z_1)G(z_{24}) + A_{4,1}^{(1)}(z_4, z_3, z_2, z_1)G(z_{34}) \\ &+ A_{4,4}^{(1)}(z_1, z_2, z_3, z_4)G(z_{12})G(z_{34}) + A_{4,4}^{(1)}(z_1, z_3, z_2, z_4)G(z_{13})G(z_{24}) \\ &+ A_{4,5}^{(1)}(z_1, z_4, z_3, z_2)G(z_{14})G(z_{23}), \\ A_{4,i}^{(1)}(z_1, z_2, z_3, z_4) &= \frac{Q_{4,i}^{(1)}(z_1, z_2, z_3, z_4)}{z_{13}z_{14}z_{23}z_{24}}, \quad (i = 1, \dots, 5), \end{aligned} \quad (4.9)$$

where we have taken the symmetry relations (4.3) into account. We can assume the polynomials  $Q_{4,i}^{(1)}(z_1, z_2, z_3, z_4)$  as

$$Q_{4,i}^{(1)}(z_1, z_2, z_3, z_4) = \sum_{\substack{0 \leq j_1, j_2, j_3, j_4 \leq 2 \\ j_1 + j_2 + j_3 + j_4 \leq 4}} C_{4,i}(j_1, j_2, j_3, j_4) z_1^{j_1} z_2^{j_2} z_3^{j_3} z_4^{j_4}. \quad (4.10)$$

Substituting (4.9) and (4.10) into (4.4), ..., (4.8), we can extract the equations for coefficients  $C_{4,i}(j_1, j_2, j_3, j_4)$  and solve them. Actually we have made a *Mathematica* program to perform the calculation. The

result is

$$\begin{aligned}
Q_{4,1}^{(1)}(z_1, z_2, z_3, z_4) &= \frac{19}{90} + \frac{7}{80}z_{12}^2 - \frac{1}{12}z_{13}^2 - \frac{1}{144}z_{14}^2 - \frac{1}{36}z_{23}^2 - \frac{1}{48}z_{24}^2 + \frac{1}{30}z_{34}^2 - \frac{1}{144}z_{13}^2z_{24}^2 + \frac{1}{144}z_{14}^2z_{23}^2 + \frac{1}{720}z_{12}^2z_{34}^2 \\
&+ i\left(\frac{1}{9}z_{13} - \frac{1}{36}z_{24} + \frac{1}{72}z_{12}^2z_{34} - \frac{1}{18}z_{12}z_{34}^2 - \frac{1}{36}z_{13}^2z_{24} + \frac{5}{72}z_{13}z_{24}^2 - \frac{1}{24}z_{14}^2z_{23}\right), \\
Q_{4,2}^{(1)}(z_1, z_2, z_3, z_4) &= \frac{19}{90} + \frac{59}{360}z_{12}^2 - \frac{1}{12}z_{13}^2 - \frac{1}{12}z_{23}^2 - \frac{1}{12}z_{14}^2 - \frac{1}{12}z_{24}^2 + \frac{4}{45}z_{34}^2 + \frac{1}{18}z_{13}^2z_{24}^2 + \frac{1}{18}z_{14}^2z_{23}^2 - \frac{17}{360}z_{12}^2z_{34}^2 \\
&+ i\left(\frac{1}{9}z_{13} + \frac{1}{9}z_{24} - \frac{1}{18}z_{13}^2z_{24} - \frac{1}{18}z_{13}z_{24}^2 - \frac{1}{18}z_{14}^2z_{23} - \frac{1}{18}z_{14}z_{23}^2\right), \\
Q_{4,3}^{(1)}(z_1, z_2, z_3, z_4) &= \frac{19}{90} + \frac{29}{360}z_{12}^2 - \frac{1}{48}z_{13}^2 - \frac{1}{48}z_{23}^2 - \frac{1}{48}z_{14}^2 - \frac{1}{48}z_{24}^2 + \frac{17}{360}z_{34}^2 + \frac{1}{72}z_{13}^2z_{24}^2 + \frac{1}{72}z_{14}^2z_{23}^2 - \frac{1}{180}z_{12}^2z_{34}^2 \\
&+ i\left(-\frac{1}{36}z_{13} - \frac{1}{36}z_{24} + \frac{1}{72}z_{13}^2z_{24} + \frac{1}{72}z_{13}z_{24}^2 + \frac{1}{72}z_{14}^2z_{23} + \frac{1}{72}z_{14}z_{23}^2\right), \\
Q_{4,4}^{(1)}(z_1, z_2, z_3, z_4) &= \frac{7}{45} + \frac{1}{15}z_{12}^2 - \frac{1}{24}z_{13}^2 - \frac{1}{72}z_{14}^2 - \frac{1}{72}z_{23}^2 - \frac{1}{24}z_{24}^2 + \frac{1}{15}z_{34}^2 + \frac{1}{360}z_{12}^2z_{34}^2 - \frac{1}{72}z_{13}^2z_{24}^2 + \frac{1}{72}z_{14}^2z_{23}^2 \\
&+ i\left(\frac{1}{18}z_{13} - \frac{1}{18}z_{24} + \frac{1}{36}z_{12}^2z_{34} - \frac{1}{36}z_{12}z_{34}^2 - \frac{1}{18}z_{13}^2z_{24} + \frac{1}{18}z_{13}z_{24}^2\right), \\
Q_{4,5}^{(1)}(z_1, z_2, z_3, z_4) &= \frac{7}{45} + \frac{17}{180}z_{12}^2 - \frac{1}{24}z_{13}^2 - \frac{1}{24}z_{14}^2 - \frac{1}{24}z_{23}^2 - \frac{1}{24}z_{24}^2 + \frac{17}{180}z_{34}^2 - \frac{1}{90}z_{12}^2z_{34}^2 + \frac{1}{36}z_{13}^2z_{24}^2 + \frac{1}{36}z_{14}^2z_{23}^2 \\
&+ i\left(\frac{1}{18}z_{13} + \frac{1}{18}z_{24} - \frac{1}{36}z_{13}^2z_{24} - \frac{1}{36}z_{13}z_{24}^2 - \frac{1}{36}z_{14}^2z_{23} - \frac{1}{36}z_{14}z_{23}^2\right). \tag{4.11}
\end{aligned}$$

Unfortunately we could not find a further compact way to express the results above. Other correlation functions in this sector are calculated by the intertwining relations. As it is not possible to present all the explicit results here, we instead show the formal expressions in Table 2. Here  $C(z_1, z_2, z_3, z_4)$  and  $Z(z_1, z_2, z_3, z_4)$  are respectively the real and the imaginary part of the correlation function  $P_{+++-}^{--+}(z_1, z_2, z_3, z_4)$  (4.11). Similarly,  $A(z_1, z_2, z_3, z_4)$ ,  $B(z_1, z_2, z_3, z_4)$ ,  $D(z_1, z_2, z_3, z_4)$ ,  $E(z_1, z_2, z_3, z_4)$  and  $X(z_1, z_2, z_3, z_4)$ ,  $Y(z_1, z_2, z_3, z_4)$  are respectively the real and the imaginary parts of other correlation functions. Using the intertwining relations, they are

$P_{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4}^{\epsilon'_1 \epsilon'_2 \epsilon'_3 \epsilon'_4}$	$- + ++$	$+ - ++$	$+ + - +$	$+ + + -$
$- + ++$	$P_3(z_2, z_3, z_4)$ $- P_4(z_1, z_2, z_3, z_4)$	$A(z_1, z_2, z_3, z_4)$ $- iX(z_1, z_2, z_3, z_4)$	$B(z_1, z_2, z_3, z_4)$ $- iY(z_1, z_2, z_3, z_4)$	$C(z_1, z_2, z_3, z_4)$ $- iZ(z_1, z_2, z_3, z_4)$
$+ - ++$	$A(z_1, z_2, z_3, z_4)$ $+ iX(z_1, z_2, z_3, z_4)$	$D(z_1, z_2, z_3, z_4)$	$E(z_1, z_2, z_3, z_4)$ $+ \frac{i}{24} (G_{123}^A + G_{234}^A)$	$B(z_4, z_3, z_2, z_1)$ $- iY(z_4, z_3, z_2, z_1)$
$+ + - +$	$B(z_1, z_2, z_3, z_4)$ $+ iY(z_1, z_2, z_3, z_4)$	$E(z_1, z_2, z_3, z_4)$ $- \frac{i}{24} (G_{123}^A + G_{234}^A)$	$D(z_4, z_3, z_2, z_1)$	$A(z_4, z_3, z_2, z_1)$ $- iX(z_4, z_3, z_2, z_1)$
$+ + + -$	$C(z_1, z_2, z_3, z_4)$ $+ iZ(z_1, z_2, z_3, z_4)$	$B(z_4, z_3, z_2, z_1)$ $+ iY(z_4, z_3, z_2, z_1)$	$A(z_4, z_3, z_2, z_1)$ $+ iX(z_4, z_3, z_2, z_1)$	$P_3(z_1, z_2, z_3)$ $- P_4(z_1, z_2, z_3, z_4)$

Table 2: Inhomogeneous correlation functions for  $n = 4$  in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 + \epsilon_4 = 2$

successively calculated from  $C(z_1, z_2, z_3, z_4)$  and  $Z(z_1, z_2, z_3, z_4)$ ,

$$\begin{aligned}
B(z_1, z_2, z_3, z_4) &= C(z_1, z_2, z_4, z_3) + \frac{1}{z_{34}} \{Z(z_1, z_2, z_3, z_4) - Z(z_1, z_2, z_4, z_3)\}, \\
Y(z_1, z_2, z_3, z_4) &= Z(z_1, z_2, z_4, z_3) + \frac{1}{z_{34}} \{C(z_1, z_2, z_4, z_3) - C(z_1, z_2, z_3, z_4)\}, \\
A(z_1, z_2, z_3, z_4) &= B(z_1, z_3, z_2, z_4) + \frac{1}{z_{23}} \{Y(z_1, z_2, z_3, z_4) - Y(z_1, z_3, z_2, z_4)\}, \\
X(z_1, z_2, z_3, z_4) &= Y(z_1, z_3, z_2, z_4) + \frac{1}{z_{23}} \{B(z_1, z_3, z_2, z_4) - B(z_1, z_2, z_3, z_4)\}, \\
D(z_1, z_2, z_3, z_4) &= P_3(z_1, z_3, z_4) - P_4(z_2, z_1, z_3, z_4) + \frac{1}{z_{21}} \{X(z_1, z_2, z_3, z_4) - X(z_2, z_1, z_3, z_4)\}, \\
E(z_1, z_2, z_3, z_4) &= B(z_2, z_1, z_3, z_4) + \frac{1}{z_{21}} \{Y(z_1, z_2, z_3, z_4) - Y(z_2, z_1, z_3, z_4)\}. \tag{4.12}
\end{aligned}$$

Note that the diagonal correlation functions  $P_{+-++}^{+-++}(z_1, z_2, z_3, z_4)$  and  $P_{++-+}^{++-+}(z_1, z_2, z_3, z_4)$  have no imaginary parts. Note also that we have found the imaginary part of  $P_{++-+}^{+-++}(z_1, z_2, z_3, z_4)$  is simply expressed in terms of the function  $G_{123}^A$  introduced in (3.1). These properties are reflected in some symmetry relations of the functions  $A(z_1, z_2, z_3, z_4), \dots, Z(z_1, z_2, z_3, z_4)$ . For example we have

$$\begin{aligned}
A(z_1, z_2, z_3, z_4) &= A(z_2, z_1, z_3, z_4), \quad E(z_1, z_2, z_3, z_4) = E(z_1, z_3, z_2, z_4), \\
Y(z_1, z_2, z_3, z_4) &= Y(z_2, z_1, z_3, z_4) + \frac{1}{24} (G_{123}^A + G_{234}^A - G_{213}^A - G_{134}^A). \tag{4.13}
\end{aligned}$$

$P_{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4}^{\epsilon'_1 \epsilon'_2 \epsilon'_3 \epsilon'_4}$	$--++$	$-+-+$	$-++-$	$+--+$	$+--+$	$++--$
$--++$	$C_1(\{z_i\})$	$C_2(\{-z_i\})$	$C_3(\{-z_i\})$	$C_4(\{-z_i\})$	$C_5(\{-z_i\})$	$C_6(\{z_i\})$
$-+-+$	$C_2(\{z_i\})$	$C_7(\{z_i\})$	$C_8(\{-z_i\})$	$C_9(\{-z_i\})$	$C_{10}(\{z_i\})$	$C_5(\{z_i\})$
$-++-$	$C_3(\{z_i\})$	$C_8(\{z_i\})$	$C_{11}(\{z_i\})$	$C_{12}(\{z_i\})$	$C_9(\{z_i\})$	$C_4(\{z_i\})$
$+--+$	$C_4(\{z_i\})$	$C_9(\{z_i\})$	$C_{12}(\{z_i\})$	$C_{11}(\{z_i\})$	$C_8(\{z_i\})$	$C_3(\{z_i\})$
$+--+$	$C_5(\{z_i\})$	$C_{10}(\{z_i\})$	$C_9(\{-z_i\})$	$C_8(\{-z_i\})$	$C_7(\{z_i\})$	$C_2(\{z_i\})$
$++--$	$C_6(\{z_i\})$	$C_5(\{-z_i\})$	$C_4(\{-z_i\})$	$C_3(\{-z_i\})$	$C_2(\{-z_i\})$	$C_1(\{z_i\})$

Table 3: Inhomogeneous correlation functions for  $n = 4$  in the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 + \epsilon_4 = 0$

We also remark the following relations hold from the identity relations.

$$\begin{aligned}
 A(z_1, z_2, z_3, z_4) + B(z_1, z_1, z_3, z_4) + C(z_1, z_2, z_3, z_4) &= 2P_4(z_1, z_2, z_3, z_4) - P_3(z_2, z_3, z_4), \\
 A(z_1, z_2, z_3, z_4) + D(z_1, z_2, z_3, z_4) + E(z_1, z_2, z_3, z_4) + B(z_4, z_3, z_2, z_1) &= P_4(z_1, z_2, z_3, z_4), \\
 X(z_1, z_2, z_3, z_4) - Y(z_4, z_3, z_2, z_1) - \frac{1}{24}(G_{123}^A + G_{234}^A) &= 0, \\
 X(z_1, z_2, z_3, z_4) + Y(z_1, z_2, z_3, z_4) + Z(z_1, z_2, z_3, z_4) &= 0. \tag{4.14}
 \end{aligned}$$

Unfortunately we have not been able to afford a nicer compact way to describe the functions  $A(z_1, z_2, z_3, z_4), \dots, Z(z_1, z_2, z_3, z_4)$ , from which the relations such as (4.13), (4.14) follow naturally.

Next let us consider the sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 + \epsilon_4 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 + \epsilon'_4 = 0$ . There are 36 correlation functions in this sector. But according to (2.8), we need to know only 12 correlation functions  $C_1(z_1, z_2, z_3, z_4), \dots, C_{12}(z_1, z_2, z_3, z_4)$  as shown in Table 3.

In the present case, they can be derived from those we have obtained already as follows.

$$\begin{aligned}
C_1(z_1, z_2, z_3, z_4) &= P_{-++}^{---}(z_2, z_3, z_4) - P_{+-+}^{---}(z_1, z_2, z_3, z_4), \\
&= \frac{1}{12} + \frac{1}{12}(G_{34} - G_{23} - G_{24}) + \frac{1}{12}G_{234} - D(z_1, z_2, z_3, z_4), \\
C_2(z_1, z_2, z_3, z_4) &= P_{+-+}^{---}(z_2, z_3, z_4) - P_{++-}^{---}(z_1, z_2, z_3, z_4), \\
&= \frac{1}{12} + \frac{1}{6}G_{12} - E(z_1, z_2, z_3, z_4) + \frac{i}{24}(G_{123}^A - G_{234}^A), \\
C_3(z_1, z_2, z_3, z_4) &= P_{++-}^{---}(z_2, z_3, z_4) - P_{++-}^{---}(z_1, z_2, z_3, z_4) \\
&= \frac{1}{12} + \frac{1}{6}G_{24} - \frac{1}{6}G_{234} - B(z_4, z_3, z_2, z_1) - i\left\{Y(z_4, z_3, z_2, z_1) - \frac{1}{12}G_{234}^A\right\}, \\
C_4(z_1, z_2, z_3, z_4) &= C_3(-z_4, -z_3, -z_2, -z_1) \\
&= \frac{1}{12} + \frac{1}{6}G_{13} - \frac{1}{6}G_{123} - B(z_1, z_2, z_3, z_4) + i\left\{Y(z_1, z_2, z_3, z_4) - \frac{1}{12}G_{123}^A\right\}, \\
C_5(z_1, z_2, z_3, z_4) &= \left(1 + \frac{i}{z_{12}}\right)C_3(z_2, z_1, z_3, z_4) - \frac{i}{z_{12}}C_3(z_1, z_2, z_3, z_4), \\
C_6(z_1, z_2, z_3, z_4) &= P_4(z_1, z_2, z_3, z_4) - \sum_{j=1}^5 C_j(z_1, z_2, z_3, z_4), \\
C_7(z_1, z_2, z_3, z_4) &= P_{-+-}^{---}(z_1, z_2, z_3) - P_{-+-}^{---}(z_1, z_2, z_3, z_4) \\
&= P_{+-+}^{---}(z_1, z_2, z_3) - P_{+-+}^{---}(z_1, z_2, z_3, z_4) \\
&= \frac{1}{12} + \frac{1}{12}(G_{13} - G_{23} - G_{12}) - \frac{1}{12}G_{123} - D(z_1, z_2, z_3, z_4), \\
C_8(z_1, z_2, z_3, z_4) &= P_{++-}^{---}(z_2, z_3, z_4) - P_{++-}^{---}(z_1, z_2, z_3, z_4) \\
&= \frac{1}{12} + \frac{1}{6}G_{34} - A(z_4, z_3, z_2, z_1) - i\left\{X(z_4, z_3, z_2, z_1) + \frac{1}{12}G_{234}^A\right\}, \\
C_9(z_1, z_2, z_3, z_4) &= C_8(-z_4, -z_3, -z_2, -z_1) \\
&= \frac{1}{12} + \frac{1}{6}G_{12} - A(z_1, z_2, z_3, z_4) + i\left\{X(z_1, z_2, z_3, z_4) + \frac{1}{12}G_{123}^A\right\}, \\
C_{10}(z_1, z_2, z_3, z_4) &= C_6(z_1, z_3, z_2, z_4) + \frac{i}{z_{32}}\{C_5(z_1, z_3, z_2, z_4) - C_5(z_1, z_2, z_3, z_4)\}, \\
C_{11}(z_1, z_2, z_3, z_4) &= P_{-++}^{---}(z_1, z_2, z_3) - P_{-++}^{---}(z_1, z_2, z_3, z_4) \\
&= -\frac{1}{6} - \frac{1}{12}(G_{12} + G_{13} + G_{24} + G_{34}) + \frac{1}{12}(G_{123} + G_{234}) + P_4(z_1, z_2, z_3, z_4), \\
C_{12}(z_1, z_2, z_3, z_4) &= C_{10}(z_1, z_2, z_4, z_3) + \frac{i}{z_{43}}\{C_9(z_1, z_2, z_4, z_3) - C_9(z_1, z_2, z_3, z_4)\}. \tag{4.15}
\end{aligned}$$

Among the correlation functions above, let us pick up a diagonal one  $P_{--++}^{--+}(z_1, z_2, z_3, z_4)$ , since it has a relatively simple explicit form.

$$\begin{aligned} P_{--++}^{--+}(z_1, z_2, z_3, z_4) &= \frac{1}{30} + A_{4,1}^{(2)}(z_1, z_2, z_3, z_4)G(z_{12}) + \left\{ A_{4,1}^{(2)}(z_1, z_3, z_2, z_4) - \frac{1}{12} \right\} G(z_{13}) \\ &\quad + \left\{ A_{4,1}^{(2)}(z_1, z_4, z_3, z_2) - \frac{1}{12} \right\} G(z_{14}) + \left\{ A_{4,1}^{(2)}(z_3, z_2, z_1, z_4) - \frac{1}{12} \right\} G(z_{23}) \\ &\quad + \left\{ A_{4,1}^{(2)}(z_4, z_2, z_3, z_1) - \frac{1}{12} \right\} G(z_{24}) + A_{4,1}^{(2)}(z_4, z_3, z_2, z_1)G(z_{34}) \\ &\quad + A_{4,2}^{(2)}(z_1, z_2, z_3, z_4)G(z_{12})G(z_{34}) + A_{4,2}^{(2)}(z_1, z_3, z_2, z_4)G(z_{13})G(z_{24}) \\ &\quad + A_{4,2}^{(2)}(z_1, z_4, z_3, z_2)G(z_{14})G(z_{23}), \end{aligned}$$

with

$$\begin{aligned} A_{4,i}^{(2)}(z_1, z_2, z_3, z_4) &= \frac{Q_{4,i}^{(2)}(z_1, z_2, z_3, z_4)}{z_{13}z_{14}z_{23}z_{24}}, \quad (i = 1, 2), \\ Q_{4,2}^{(2)}(z_1, z_2, z_3, z_4) &= \frac{1}{36} \left\{ (z_{13}z_{24} + 1)(z_{14}z_{23} + 1) + \frac{2}{5} \left( z_{12}^2 + \frac{3}{2} \right) \left( z_{34}^2 + \frac{3}{2} \right) - \frac{5}{2} \right\}, \\ Q_{4,1}^{(2)}(z_1, z_2, z_3, z_4) &= 2Q_{4,2}^{(2)}(z_1, z_2, z_3, z_4) - \frac{1}{60}(z_{12}^2 - 1)(z_{34}^2 + 1). \end{aligned} \quad (4.16)$$

The correlation functions in the remaining sector  $\epsilon_1 + \epsilon_2 + \epsilon_3 + \epsilon_4 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 + \epsilon'_4 = -2$  are obtained from those in  $\epsilon_1 + \epsilon_2 + \epsilon_3 + \epsilon_4 = \epsilon'_1 + \epsilon'_2 + \epsilon'_3 + \epsilon'_4 = 2$  by the negation relation

$$P_{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4}^{\epsilon'_1 \epsilon'_2 \epsilon'_3 \epsilon'_4}(z_1, z_2, z_3, z_4) = P_{-\epsilon_1, -\epsilon_2, -\epsilon_3, -\epsilon_4}^{-\epsilon'_1, -\epsilon'_2, -\epsilon'_3, -\epsilon'_4}(z_1, z_2, z_3, z_4). \quad (4.17)$$

In this way we can get all the inhomogeneous correlation functions for  $n = 4$ .

Now we shortly consider the diagonal elements of the correlation functions  $P_{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4}^{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4}(z_1, z_2, z_3, z_4)$ . We claim they can be expressed only in terms of the EFP's,  $P_2(z_1, z_2)$ ,  $P_3(z_1, z_2, z_3)$ ,  $P_4(z_1, z_2, z_3, z_4)$  and  $P_{--++}^{--+}(z_1, z_2, z_3, z_4)$ . For example, we have

$$\begin{aligned} P_{-+++}^{----}(z_1, z_2, z_3, z_4) &= P_3(z_2, z_3, z_4) - P_4(z_1, z_2, z_3, z_4), \\ P_{+-++}^{----}(z_1, z_2, z_3, z_4) &= P_2(z_3, z_4) - P_3(z_2, z_3, z_4) - P_{--++}^{--+}(z_1, z_2, z_3, z_4), \\ P_{+-+-}^{----}(z_1, z_2, z_3, z_4) &= \frac{1}{2} - P_2(z_1, z_2) - P_2(z_2, z_3) - P_2(z_3, z_4) \\ &\quad + P_3(z_1, z_2, z_3) + P_3(z_2, z_3, z_4) + P_{--++}^{--+}(z_1, z_2, z_3, z_4), \\ P_{+--+}^{----}(z_1, z_2, z_3, z_4) &= P_2(z_2, z_3) - P_3(z_1, z_2, z_3) - P_3(z_2, z_3, z_4) + P_4(z_1, z_2, z_3, z_4). \end{aligned} \quad (4.18)$$

Then, especially, the inhomogeneous generalization of the third-neighbor correlator  $\langle S_1^z S_4^z \rangle$  is represented

$$\begin{aligned}
& \langle S_1^z S_4^z \rangle(z_1, z_2, z_3, z_4) \\
&= \frac{1}{4} \left\{ P_{+++\bar{+}}^{++++}(z_1, z_2, z_3, z_4) - P_{-++\bar{+}}^{-+++}(z_1, z_2, z_3, z_4) + P_{+-+\bar{+}}^{+-++}(z_1, z_2, z_3, z_4) - P_{--+\bar{+}}^{--++}(z_1, z_2, z_3, z_4) \right. \\
&\quad + P_{++-\bar{+}}^{++-+}(z_1, z_2, z_3, z_4) - P_{-+-\bar{+}}^{-+-+}(z_1, z_2, z_3, z_4) + P_{+-\bar{+}}^{+-+-}(z_1, z_2, z_3, z_4) - P_{--\bar{+}}^{--+-}(z_1, z_2, z_3, z_4) \\
&\quad - P_{+++\bar{-}}^{+++-}(z_1, z_2, z_3, z_4) + P_{-++\bar{-}}^{-++-}(z_1, z_2, z_3, z_4) - P_{+-+\bar{-}}^{+-+-}(z_1, z_2, z_3, z_4) + P_{--+\bar{-}}^{--+-}(z_1, z_2, z_3, z_4) \\
&\quad - P_{++-\bar{-}}^{++--}(z_1, z_2, z_3, z_4) + P_{-+-\bar{-}}^{-+-+}(z_1, z_2, z_3, z_4) - P_{+-\bar{-}}^{+---}(z_1, z_2, z_3, z_4) + P_{--\bar{-}}^{----}(z_1, z_2, z_3, z_4) \Big\} \\
&= -\frac{1}{4} + P_2(z_1, z_2) + P_2(z_2, z_3) + P_2(z_3, z_4) - 2 \{P_3(z_1, z_2, z_3) + P_3(z_2, z_3, z_4)\} \\
&\quad + 2 \{P_4(z_1, z_2, z_3, z_4) - P_{--+\bar{+}}^{--++}(z_1, z_2, z_3, z_4)\} \\
&= \frac{1}{12} + \frac{1}{6} G_{14} + \frac{1}{6} (G_{123} + G_{234}) + A_{4,1}^{(3)}(z_1, z_2, z_3, z_4) G_{12} + A_{4,1}^{(3)}(z_1, z_3, z_2, z_4) G_{13} \\
&\quad + A_{4,1}^{(3)}(z_1, z_4, z_3, z_2) G_{14} + A_{4,1}^{(3)}(z_3, z_2, z_1, z_4) G_{23} + A_{4,1}^{(3)}(z_4, z_2, z_3, z_1) G_{24} \\
&\quad + A_{4,1}^{(3)}(z_4, z_3, z_2, z_1) G_{34} + A_{4,2}^{(3)}(z_1, z_2, z_3, z_4) G_{12} G_{34} + A_{4,2}^{(3)}(z_1, z_3, z_2, z_4) G_{13} G_{24} \\
&\quad + A_{4,2}^{(3)}(z_1, z_4, z_3, z_2) G_{14} G_{23},
\end{aligned}$$

where

$$\begin{aligned}
A_{4,2}^{(3)}(z_1, z_2, z_3, z_4) &= 2 \left\{ A_{4,2}^{(0)}(z_1, z_2, z_3, z_4) - A_{4,2}^{(3)}(z_1, z_2, z_3, z_4) \right\} \\
&= \frac{1}{9} \left( \frac{2}{z_{13} z_{14} z_{23} z_{24}} - \frac{1}{z_{14} z_{23}} - \frac{1}{z_{13} z_{24}} \right), \\
A_{4,1}^{(3)}(z_1, z_2, z_3, z_4) &= 2 \left\{ A_{4,1}^{(0)}(z_1, z_2, z_3, z_4) - A_{4,1}^{(3)}(z_1, z_2, z_3, z_4) \right\} \\
&= \frac{1}{18} \left( \frac{5}{z_{13} z_{14} z_{23} z_{24}} - \frac{1}{z_{14} z_{23}} - \frac{1}{z_{13} z_{24}} - \frac{3}{z_{14} z_{24}} - \frac{3}{z_{13} z_{23}} \right). \quad (4.19)
\end{aligned}$$

By taking the homogeneous limit  $z_i \rightarrow 0$ , we recover the result (1.6).

$$\begin{aligned}
\lim_{z_i \rightarrow 0} \langle S_1^z S_4^z \rangle(z_1, z_2, z_3, z_4) &= \frac{1}{12} - \frac{1}{3} \zeta_a(1) + 2 \{\zeta_a(1) - \zeta_a(3)\} \\
&\quad + \left\{ -\frac{14}{3} \zeta_a(1) + \frac{92}{9} \zeta_a(3) - \frac{50}{9} \zeta_a(5) - \frac{56}{9} \zeta_a(1) \zeta_a(3) - \frac{8}{3} \zeta_a(3)^2 + \frac{80}{9} \zeta_a(1) \zeta_a(5) \right\} \\
&= \frac{1}{12} - 3 \zeta_a(1) + \frac{74}{9} \zeta_a(3) - \frac{50}{9} \zeta_a(5) - \frac{56}{9} \zeta_a(1) \zeta_a(3) - \frac{8}{3} \zeta_a(3)^2 + \frac{80}{9} \zeta_a(1) \zeta_a(5) \\
&= \frac{1}{12} - 3 \ln 2 + \frac{37}{6} \zeta(3) - \frac{125}{24} \zeta(5) - \frac{14}{3} \ln 2 \cdot \zeta(3) - \frac{3}{2} \zeta(3)^2 + \frac{25}{3} \ln 2 \cdot \zeta(5). \quad (4.20)
\end{aligned}$$

Other correlation functions in (1.6) are reproduced in a similar way.

## 5 Correlation functions in the case $n = 5$

Following the scheme in the previous sections we could further obtain all the correlation functions for the  $XXX$  chain on five lattice sites. In fact we have calculated two inhomogeneous correlation functions  $P_{++++-}^{----}(z_1, z_2, z_3, z_4, z_5)$  and  $P_{+++-}^{--++}(z_1, z_2, z_3, z_4, z_5)$  from the first principle relations. Other correlation functions are derived by the intertwining relations, etc..., as before. The calculations have been performed by the heavy use of *Mathematica* and the obtained results are too complicated to be described thoroughly in this paper. Therefore here we shall omit to describe them. We, however, give explicit form of the inhomogeneous generalization of the fourth-neighbor correlation function  $\langle S_1^z S_5^z \rangle(z_1, z_2, z_3, z_4, z_5)$ .

$$\begin{aligned} \langle S_1^z S_5^z \rangle(z_1, z_2, z_3, z_4, z_5) &\equiv \frac{1}{4} \sum_{\epsilon_j=\pm} \epsilon_1 \epsilon_5 P_{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4 \epsilon_5}^{\epsilon_1 \epsilon_2 \epsilon_3 \epsilon_4 \epsilon_5}(z_1, z_2, z_3, z_4, z_5) \\ &= \frac{1}{12} + A_{5,1}(z_1, z_2, z_3, z_4, z_5)G_{12} + A_{5,1}(z_1, z_3, z_2, z_4, z_5)G_{13} + A_{5,1}(z_1, z_4, z_3, z_2, z_5)G_{14} \\ &+ A_{5,1}(z_5, z_2, z_3, z_4, z_1)G_{25} + A_{5,1}(z_5, z_3, z_2, z_4, z_1)G_{35} + A_{5,1}(z_5, z_4, z_3, z_2, z_1)G_{45} \\ &+ A_{5,2}(z_2, z_3, z_1, z_4, z_5)G_{23} + A_{5,2}(z_2, z_4, z_1, z_3, z_5)G_{24} + A_{5,2}(z_3, z_4, z_1, z_2, z_5)G_{34} \\ &+ A_{5,3}(z_1, z_5, z_2, z_3, z_4)G_{15} + A_{5,4}(z_1, z_2, z_3, z_4, z_5)G_{12}G_{34} + A_{5,4}(z_1, z_3, z_2, z_4, z_5)G_{13}G_{24} \\ &+ A_{5,4}(z_1, z_4, z_2, z_3, z_5)G_{14}G_{23} + A_{5,4}(z_5, z_2, z_3, z_4, z_1)G_{25}G_{34} + A_{5,4}(z_5, z_3, z_2, z_4, z_1)G_{35}G_{24} \\ &+ A_{5,4}(z_5, z_4, z_2, z_3, z_1)G_{45}G_{23} + A_{5,5}(z_1, z_2, z_3, z_5, z_4)G_{12}G_{35} + A_{5,5}(z_1, z_3, z_2, z_5, z_4)G_{13}G_{25} \\ &+ A_{5,4}(z_1, z_2, z_4, z_5, z_3)G_{12}G_{45} + A_{5,5}(z_1, z_4, z_2, z_5, z_3)G_{14}G_{25} + A_{5,5}(z_1, z_3, z_4, z_5, z_2)G_{13}G_{45} \\ &+ A_{5,4}(z_1, z_4, z_3, z_5, z_2)G_{14}G_{35} + A_{5,5}(z_1, z_5, z_2, z_3, z_4)G_{15}G_{23} + A_{5,5}(z_1, z_5, z_3, z_4, z_2)G_{15}G_{34} \\ &+ A_{5,5}(z_1, z_5, z_2, z_4, z_3)G_{15}G_{24}, \end{aligned}$$

where

$$A_{5,i}(z_1, z_2, z_3, z_4, z_5) = \begin{cases} \frac{Q_{5,i}(z_1, z_2, z_3, z_4, z_5)}{z_{13}z_{14}z_{15}z_{23}z_{24}z_{25}}, & (i = 1, 2, 3) \\ \frac{Q_{5,i}(z_1, z_2, z_3, z_4, z_5)}{z_{13}z_{14}z_{15}z_{23}z_{24}z_{25}z_{34}z_{35}}, & (i = 4, 5) \end{cases}$$

with

$$\begin{aligned}
Q_{5,1}(z_1, z_2, z_3, z_4, z_5) &= \frac{1}{36} [(z_{15} + z_{25})(z_{13} + z_{24})(2 - z_{13}z_{14} - z_{14}z_{23}) - 6(z_{13}z_{14} - 2)(z_{23}z_{24} - 2) \\
&\quad + z_{12}^2(z_{13}z_{23} + z_{14}z_{24} - 16) + 10z_{34}^2 - 2], \\
Q_{5,2}(z_1, z_2, z_3, z_4, z_5) &= Q_{5,1}(z_1, z_2, z_3, z_4, z_5) + \frac{1}{18}z_{13}z_{23}[5z_{14}z_{24} + z_{45}(z_{14} + z_{24}) - 5], \\
Q_{5,3}(z_1, z_2, z_3, z_4, z_5) &= Q_{5,1}(z_1, z_2, z_3, z_4, z_5) + \frac{1}{18}z_{15}z_{25}[3z_{13}z_{14}z_{23}z_{24} - 4(z_{13}z_{24} + z_{14}z_{23}) - 3z_{34}^2 + 5], \\
Q_{5,4}(z_1, z_2, z_3, z_4, z_5) &= \frac{1}{18}[2(z_{13}z_{24} + z_{14}z_{23})z_{15}z_{25} - z_{13}z_{23}z_{24}z_{45} - z_{14}z_{23}z_{24}z_{35} - z_{13}z_{14}z_{23}z_{45} \\
&\quad - z_{13}z_{14}z_{24}z_{35} + z_{12}^2z_{34}^2 + 6(z_{12}^2 + z_{34}^2) - \frac{3}{2}(z_{13}^2 + z_{14}^2 + z_{23}^2 + z_{24}^2) \\
&\quad - (z_{15}^2 + z_{25}^2 + z_{35}^2 + z_{45}^2) + 12], \\
Q_{5,5}(z_1, z_2, z_3, z_4, z_5) &= Q_{5,4}(z_1, z_2, z_3, z_4, z_5) + \frac{1}{9}(2 - z_{13}z_{24} - z_{14}z_{23})z_{15}z_{25}z_{35}z_{45}. \tag{5.21}
\end{aligned}$$

By taking the homogeneous limit, we obtain the fourth-neighbor correlation function for the homogeneous Heisenberg model as follows.

$$\begin{aligned}
\langle S_j^z S_{j+4}^z \rangle &= \lim_{z_i \rightarrow 0} \langle S_1^z S_5^z \rangle (z_1, z_2, z_3, z_4, z_5) \\
&= \frac{1}{12} - \frac{16}{3}\zeta_a(1) + \frac{290}{9}\zeta_a(3) - 72\zeta_a(1) \cdot \zeta_a(3) - \frac{1172}{9}\zeta_a(3)^2 - \frac{700}{9}\zeta_a(5) \\
&\quad + \frac{4640}{9}\zeta_a(1) \cdot \zeta_a(5) - \frac{220}{9}\zeta_a(3) \cdot \zeta_a(5) - \frac{400}{3}\zeta_a(5)^2 + \frac{455}{9}\zeta_a(7) \\
&\quad - \frac{3920}{9}\zeta_a(1) \cdot \zeta_a(7) + 280\zeta_a(3) \cdot \zeta_a(7) \\
&= \frac{1}{12} - \frac{16}{3}\ln 2 + \frac{145}{6}\zeta(3) - 54\ln 2 \cdot \zeta(3) - \frac{293}{4}\zeta(3)^2 - \frac{875}{12}\zeta(5) \\
&\quad + \frac{1450}{3}\ln 2 \cdot \zeta(5) - \frac{275}{16}\zeta(3) \cdot \zeta(5) - \frac{1875}{16}\zeta(5)^2 + \frac{3185}{64}\zeta(7) \\
&\quad - \frac{1715}{4}\ln 2 \cdot \zeta(7) + \frac{6615}{32}\zeta(3) \cdot \zeta(7) = 0.034652776982 \dots. \tag{5.22}
\end{aligned}$$

This is one of the main new results of this paper. Furthermore we have obtained all the other correlation functions for  $n = 5$ . We present below the homogeneous limit of independent ones.

$$\begin{aligned}
\langle S_j^x S_{j+1}^x S_{j+2}^z S_{j+4}^z \rangle &= \frac{1}{240} + \frac{1}{12}\ln 2 - \frac{517}{480}\zeta(3) + \frac{25}{12}\ln 2 \cdot \zeta(3) + \frac{203}{80}\zeta(3)^2 + \frac{1525}{384}\zeta(5) \\
&\quad - \frac{215}{12}\ln 2 \cdot \zeta(5) + \frac{5}{16}\zeta(3) \cdot \zeta(5) + \frac{125}{32}\zeta(5)^2 - \frac{735}{256}\zeta(7) \\
&\quad + \frac{1029}{64}\ln 2 \cdot \zeta(7) - \frac{441}{64}\zeta(3) \cdot \zeta(7) = -0.0098924350847 \dots, \tag{5.23}
\end{aligned}$$

$$\begin{aligned} \langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+4}^z \rangle &= \frac{1}{240} - \frac{1}{4} \ln 2 + \frac{301}{160} \zeta(3) - \frac{9}{2} \ln 2 \cdot \zeta(3) - \frac{1079}{160} \zeta(3)^2 - \frac{975}{128} \zeta(5) \\ &\quad + \frac{365}{8} \ln 2 \cdot \zeta(5) - \frac{185}{128} \zeta(3) \cdot \zeta(5) - \frac{1375}{128} \zeta(5)^2 + \frac{735}{128} \zeta(7) \\ &\quad - \frac{1323}{32} \ln 2 \cdot \zeta(7) + \frac{4851}{256} \zeta(3) \cdot \zeta(7) = 0.0027889733995 \dots, \end{aligned} \quad (5.24)$$

$$\begin{aligned} \langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+4}^x \rangle &= \frac{1}{240} - \frac{5}{12} \ln 2 + \frac{569}{240} \zeta(3) - \frac{61}{12} \ln 2 \cdot \zeta(3) - \frac{1109}{160} \zeta(3)^2 - \frac{775}{96} \zeta(5) \\ &\quad + \frac{140}{3} \ln 2 \cdot \zeta(5) - \frac{185}{128} \zeta(3) \cdot \zeta(5) - \frac{1375}{128} \zeta(5)^2 + \frac{735}{128} \zeta(7) \\ &\quad - \frac{1323}{32} \ln 2 \cdot \zeta(7) + \frac{4851}{256} \zeta(3) \cdot \zeta(7) = -0.00505666089481 \dots, \end{aligned} \quad (5.25)$$

$$\begin{aligned} \langle S_j^x S_{j+1}^x S_{j+3}^z S_{j+4}^z \rangle &= \frac{1}{240} + \frac{1}{4} \ln 2 - \frac{419}{160} \zeta(3) + \frac{35}{6} \ln 2 \cdot \zeta(3) + \frac{663}{80} \zeta(3)^2 + \frac{4445}{384} \zeta(5) \\ &\quad - \frac{115}{2} \ln 2 \cdot \zeta(5) + \frac{85}{64} \zeta(3) \cdot \zeta(5) + \frac{1625}{128} \zeta(5)^2 - \frac{2303}{256} \zeta(7) \\ &\quad + \frac{833}{16} \ln 2 \cdot \zeta(7) - \frac{5733}{256} \zeta(3) \cdot \zeta(7) = 0.01857662093837 \dots, \end{aligned} \quad (5.26)$$

$$\begin{aligned} \langle S_j^x S_{j+1}^z S_{j+3}^x S_{j+4}^z \rangle &= \frac{1}{240} - \frac{5}{12} \ln 2 + \frac{1883}{480} \zeta(3) - \frac{28}{3} \ln 2 \cdot \zeta(3) - \frac{551}{40} \zeta(3)^2 - \frac{2285}{128} \zeta(5) \\ &\quad + \frac{1135}{12} \ln 2 \cdot \zeta(5) - \frac{165}{64} \zeta(3) \cdot \zeta(5) - \frac{1375}{64} \zeta(5)^2 + \frac{3577}{256} \zeta(7) \\ &\quad - \frac{343}{4} \ln 2 \cdot \zeta(7) + \frac{4851}{128} \zeta(3) \cdot \zeta(7) = -0.00108936897291 \dots, \end{aligned} \quad (5.27)$$

$$\begin{aligned} \langle S_j^x S_{j+1}^z S_{j+3}^z S_{j+4}^x \rangle &= \frac{1}{240} - \frac{1}{2} \ln 2 + \frac{497}{120} \zeta(3) - \frac{28}{3} \ln 2 \cdot \zeta(3) - \frac{551}{40} \zeta(3)^2 - \frac{1155}{64} \zeta(5) \\ &\quad + \frac{1135}{12} \ln 2 \cdot \zeta(5) - \frac{165}{64} \zeta(3) \cdot \zeta(5) - \frac{1375}{64} \zeta(5)^2 + \frac{3577}{256} \zeta(7) \\ &\quad - \frac{343}{4} \ln 2 \cdot \zeta(7) + \frac{4851}{128} \zeta(3) \cdot \zeta(7) = 0.001573361370145 \dots. \end{aligned} \quad (5.28)$$

As a confirmation of our results, we have evaluated the numerical values of the correlation functions for the finite system size  $N$  up to 32 (Table 4). We have further applied an extrapolation to this data such as  $c_\infty + c_1/N^2 + c_2/N^4 + c_3/N^6$ , and obtained the estimated values  $c_\infty$  in the limit  $N \rightarrow \infty$ . One can clearly observe our analytical results coincide with the extrapolated values with extremely high accuracy (more than 4 digits).

Table 4: Numerical values of correlation functions for finite systems

	$N=26$	$N=28$	$N=30$	$N=32$	$c_\infty (N = \infty)$
$\langle S_j^z S_{j+4}^z \rangle$	0.0357233	0.0355713	0.0354497	0.0353508	0.0346535
$\langle S_j^x S_{j+1}^x S_{j+2}^z S_{j+4}^z \rangle$	-0.0099654	-0.0099553	-0.0099471	-0.0099405	-0.0098925
$\langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+4}^z \rangle$	0.0028572	0.0028477	0.0028400	0.0028337	0.0027890
$\langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+4}^x \rangle$	-0.0052334	-0.0052083	-0.0051882	-0.0051719	-0.0050564
$\langle S_j^x S_{j+1}^x S_{j+3}^z S_{j+4}^z \rangle$	0.0185725	0.0185731	0.0185737	0.0185741	0.0185766
$\langle S_j^x S_{j+1}^z S_{j+3}^x S_{j+4}^z \rangle$	-0.0011377	-0.0011309	-0.0011254	-0.0011209	-0.0010893
$\langle S_j^x S_{j+1}^z S_{j+3}^z S_{j+4}^x \rangle$	0.0016526	0.0016413	0.0016323	0.0016250	0.0015732

We make several comments on our results. First, the two other diagonal correlation functions are given as a sum of non-diagonal ones similarly in the case of four lattice sites,

$$\begin{aligned}
 \langle S_j^z S_{j+1}^z S_{j+2}^z S_{j+4}^z \rangle &= \langle S_j^x S_{j+1}^x S_{j+2}^z S_{j+4}^z \rangle + \langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+4}^z \rangle + \langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+4}^x \rangle \\
 &= \frac{1}{80} - \frac{7}{12} \ln 2 + \frac{127}{40} \zeta(3) - \frac{15}{2} \ln 2 \cdot \zeta(3) - \frac{891}{80} \zeta(3)^2 - \frac{375}{32} \zeta(5) \\
 &\quad + \frac{595}{8} \ln 2 \cdot \zeta(5) - \frac{165}{64} \zeta(3) \cdot \zeta(5) - \frac{1125}{64} \zeta(5)^2 + \frac{2205}{256} \zeta(7) \\
 &\quad - \frac{4263}{64} \ln 2 \cdot \zeta(7) + \frac{3969}{128} \zeta(3) \cdot \zeta(7) = -0.0121601225799 \dots, \quad (5.29)
 \end{aligned}$$

$$\begin{aligned}
 \langle S_j^z S_{j+1}^z S_{j+3}^z S_{j+4}^z \rangle &= \langle S_j^x S_{j+1}^x S_{j+3}^z S_{j+4}^z \rangle + \langle S_j^x S_{j+1}^z S_{j+3}^x S_{j+4}^z \rangle + \langle S_j^x S_{j+1}^z S_{j+3}^z S_{j+4}^x \rangle \\
 &= \frac{1}{80} - \frac{2}{3} \ln 2 + \frac{1307}{240} \zeta(3) - \frac{77}{6} \ln 2 \cdot \zeta(3) - \frac{1541}{80} \zeta(3)^2 - \frac{2335}{96} \zeta(5) \\
 &\quad + \frac{395}{3} \ln 2 \cdot \zeta(5) - \frac{245}{64} \zeta(3) \cdot \zeta(5) - \frac{3875}{128} \zeta(5)^2 + \frac{4851}{256} \zeta(7) \\
 &\quad - \frac{1911}{16} \ln 2 \cdot \zeta(7) + \frac{13671}{256} \zeta(3) \cdot \zeta(7) = 0.0190606133356 \dots. \quad (5.30)
 \end{aligned}$$

The expression for  $P(5)$  (1.5) is reproduced as

$$\begin{aligned}
P(5) &= \frac{1}{32} + \frac{1}{2} \langle S_j^z S_{j+1}^z \rangle + \frac{3}{8} \langle S_j^z S_{j+2}^z \rangle + \frac{1}{4} \langle S_j^z S_{j+3}^z \rangle + \frac{1}{8} \langle S_j^z S_{j+4}^z \rangle \\
&\quad + \langle S_j^z S_{j+1}^z S_{j+2}^z S_{j+3}^z \rangle + \langle S_j^z S_{j+1}^z S_{j+3}^z S_{j+4}^z \rangle + \frac{1}{2} \langle S_j^z S_{j+1}^z S_{j+3}^z S_{j+4}^z \rangle \\
&= \frac{1}{6} - \frac{10}{3} \ln 2 + \frac{281}{24} \zeta(3) - \frac{45}{2} \ln 2 \cdot \zeta(3) - \frac{489}{16} \zeta(3)^2 - \frac{6775}{192} \zeta(5) \\
&\quad + \frac{1225}{6} \ln 2 \cdot \zeta(5) - \frac{425}{64} \zeta(3) \cdot \zeta(5) - \frac{12125}{256} \zeta(5)^2 + \frac{6223}{256} \zeta(7) \\
&\quad - \frac{11515}{64} \ln 2 \cdot \zeta(7) + \frac{42777}{512} \zeta(3) \cdot \zeta(7) = 2.0117259 \dots \times 10^{-6}.
\end{aligned}$$

In fact, arbitrary correlation functions on five lattice sites are similarly calculated.

Finally we notice that the “next-nearest chiral correlator” simplifies drastically as

$$\begin{aligned}
\langle (\mathbf{S}_j \times \mathbf{S}_{j+1}) \cdot (\mathbf{S}_{j+3} \times \mathbf{S}_{j+4}) \rangle &= 6 (\langle S_j^x S_{j+1}^z S_{j+3}^x S_{j+4}^z \rangle - \langle S_j^x S_{j+1}^z S_{j+3}^z S_{j+4}^x \rangle) \\
&= \frac{1}{2} \ln 2 - \frac{21}{16} \zeta(3) + \frac{75}{64} \zeta(5) \\
&= -0.015976382058 \dots,
\end{aligned} \tag{5.31}$$

as well as a similar correlator

$$\begin{aligned}
\langle (\mathbf{S}_j \times \mathbf{S}_{j+1}) \cdot (\mathbf{S}_{j+2} \times \mathbf{S}_{j+4}) \rangle &= 6 (\langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+4}^z \rangle - \langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+4}^x \rangle) \\
&= \ln 2 - \frac{47}{16} \zeta(3) + \frac{7}{2} \ln 2 \cdot \zeta(3) + \frac{9}{8} \zeta(3)^2 + \frac{175}{64} \zeta(5) - \frac{25}{4} \ln 2 \cdot \zeta(5) \\
&= 0.04707380576617 \dots.
\end{aligned} \tag{5.32}$$

Recall that a similar simplification occurs in the case of “nearest chiral correlator” [34]

$$\begin{aligned}
\langle (\mathbf{S}_j \times \mathbf{S}_{j+1}) \cdot (\mathbf{S}_{j+2} \times \mathbf{S}_{j+3}) \rangle &= 6 (\langle S_j^x S_{j+1}^z S_{j+2}^x S_{j+3}^z \rangle - \langle S_j^x S_{j+1}^z S_{j+2}^z S_{j+3}^x \rangle) \\
&= \frac{1}{2} \ln 2 - \frac{3}{8} \zeta(3).
\end{aligned} \tag{5.33}$$

## 6 Summary

In this paper, we calculate correlation functions of the Heisenberg chain without magnetic field in the anti-ferromagnetic ground state using several fundamental functional relations followed from

the quantum Knizhnik-Zamolodchikov equations. This is the generalization of the method proposed for the emptiness formation probability. In fact we have calculated all the correlation functions on five lattice sites. Especially we could get the analytic formula for the fourth-neighbor correlation  $\langle S_j^z S_{j+4}^z \rangle$ . In principle we can continue to calculate  $\langle S_j^z S_{j+k}^z \rangle_{k \geq 5}$ , which will be reported in the further publications.

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## References

- [1] W. Heisenberg, Z. Phys. **49** (9-10) (1928) 619.
- [2] M. Takahashi, *Thermodynamics of One-Dimensional Solvable Models*, Cambridge University Press, Cambridge, 1999.
- [3] H. Bethe, Z. Phys. **76** (1931) 205.
- [4] L. Hulthén, Ark. Mat. Astron. Fysik **A 26** (1939) 1.
- [5] L. Faddeev, “*How algebraic Bethe Ansatz works for integrable models*”, Les HouchesSession LXIV (1995), Quantum symmetries, eds A. Connes, K. Gawedzki and J. Zinn-Justin, Elsevier Science (1998) pp. 149 - 219.
- [6] L. Faddeev and L. Takhtajan “*Spectrum and scattering of excitations in the one-dimensional isotropic Heisenberg model*”, J. Sov. Math. **24** (1984) 241-267.

- [7] R. Baxter, *Exactly Solved Models in Statistical Mechanics*, Academic Press, New York, 1982.
- [8] M. Takahashi, J. Phys. C: Solid State Phys. **10** (1977) 1289.
- [9] J. Dittrich and V. I. Inozemtsev, J. Phys. A **30** (1997) L623.
- [10] V.G. Knizhnik and A.B. Zamolodchikov, Nucl. Phys. B **247** (1984) 83.
- [11] F.A. Smirnov, *Form Factors in Completely Integrable Models of Quantum Field Theory*. Adv. Series in Math. Phys. 14, World Scientific, Singapore, 1992.
- [12] F.A. Smirnov, Int. J. Math. Phys. A7 (Suppl. 1B) (1992) 813.
- [13] I. Frenkel, N. Reshetikhin, Commun. Math. Phys. **146** (1992) 1.
- [14] M. Jimbo, K. Miki, T. Miwa, A. Nakayashiki, Phys. Lett. A **168** (1992) 256.
- [15] M. Jimbo and T. Miwa, *Algebraic Analysis of Solvable Lattice Models*, CBMS Regional Conference Series in Mathematics vol.**85**, American Mathematical Society, Providence, 1994.
- [16] V. Korepin, A. Izergin, F. Essler and D. Uglov, Phys. Lett. A **190** (1994) 182.
- [17] M. Jimbo, T. Miwa, J.Phys. A **29** (1996) 2923.
- [18] A. Nakayashiki, F.A. Smirnov, Commun. Math. Phys. **217** (2001) 623.
- [19] A. Nakayashiki, “*On the Cohomology of Theta Divisor of Hyperelliptic Jacobian*”, Contemporary mathematics, vol 309, in Integrable systems, topology and physics, M. Guest et al. ed., American Mathematical Society, Providence, 2002.
- [20] N. Kitanine, J.M. Maillet, V. Terras, Nucl. Phys. B **554** (1999), 647.
- [21] N. Kitanine, J.M. Maillet, V. Terras, Nucl. Phys. B **567** (2000), 554.
- [22] H.E. Boos, V.E. Korepin, J. Phys. A **34** (2001) 5311.

- [23] H.E. Boos, V.E. Korepin, “*Evaluation of integrals representing correlators in XXX Heisenberg spin chain*” in. MathPhys Odyssey 2001, Birkhäuser, Basel, (2001) 65.
- [24] H.E. Boos, V.E. Korepin, F.A. Smirnov, Nucl. Phys. **B** 658 (2003) 417.
- [25] H.E. Boos, V.E. Korepin and F.A. Smirnov, J. Phys. A: Math. Gen. **37** (2004) 323.
- [26] H.E. Boos, V.E. Korepin and F.A. Smirnov, “*New formulae for solutions of quantum Knizhnik-Zamolodchikov equations on level  $-4$  and correlation functions*”, hep-th/0305135, to appear in Moscow Mathematical Journal.
- [27] H.E. Boos, V.E. Korepin and F.A. Smirnov, “*Connecting lattice and relativistic models via conformal field theory*”, math-ph/0311020, Prog. Math. in press.
- [28] H.E. Boos, V.E. Korepin, Y. Nishiyama and M. Shiroishi, J. Phys. A: Math. Gen **35** (2002) 4443.
- [29] K. Sakai, M. Shiroishi, Y. Nishiyama, M. Takahashi, Phys. Rev. E **67** (2003) 065101.
- [30] G. Kato, M. Shiroishi, M. Takahashi, K. Sakai, J. Phys. A: Math. Gen. **36** (2003) L337.
- [31] M. Takahashi, G. Kato, M. Shiroishi, J. Phys. Soc. Jpn. **73** (2004) 245.
- [32] G. Kato, M. Shiroishi, M. Takahashi, K. Sakai, J. Phys. A: Math. Gen. **37** (2004) 5097.
- [33] H. Boos, M. Jimbo, T. Miwa, F. Smirnov and Y. Takeyama, “*A recursion formula for the correlation functions of an inhomogeneous XXX model*”, hep-th/0405044.
- [34] N. Muramoto and M. Takahashi, J. Phys. Soc. Jpn. **68** (1999) 2098.